



## 7th International Conference on Particle Physics and Astrophysics

Open and hidden strangeness in central A+A collisions:  
ratios of the average transverse energy density for  $\phi$  –meson  
and  $\Omega$  –hyperons in the range  $\sqrt{s_{NN}} = 39 \text{ GeV} - 2,76 \text{ TeV}$

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2.

# Interesting strangeness

## Hadrons with (multiple) strange quarks

Small hadronic cross section

Sensitive to dynamics of the medium

Can be easily reconstructed and identified in experiment!

→ Systematic study of medium properties!

### $\phi$ -meson (1020)

The lightest meson with hidden flavor. Contains only strange quarks.

→ A dramatic increase prediction in the  $\phi$  -meson after the formation of the quark-gluon plasma (QGP). [1]

[1] - Asher Shor , PHYSICAL REVIEW LETTERS, 25.70.Np, 12.35.Ht, 21.65 (1985)

3.

# Motivation

**The ratio of particles with different strangeness can indicate whether there is an increased production of  $\phi$  mesons.**

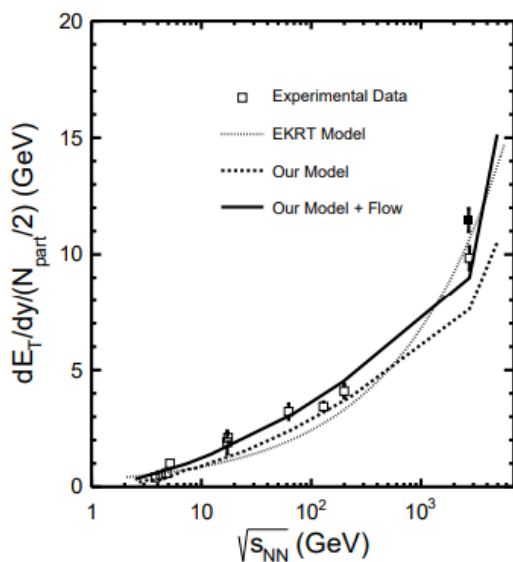
Particle	<u>кварковый состав</u>	<u>m, MeV</u>	$\tau, c$
$\phi$	$s\bar{s}$	1020	$1,54 \cdot 10^{-22}$
$\Omega$	$sss$	1672	$8,21 \cdot 10^{-11}$

$\phi, \Omega, \Omega^-$

4.

# Transverse energy

$$\langle m_{\perp} \rangle = \sqrt{m^2 + \langle p_{\perp} \rangle^2}$$



[3]

[2]

$$\frac{d \langle E_{\perp} \rangle}{dy} = \langle m_{\perp} \rangle \frac{dN}{dy}$$

$S_{\perp}$  - is the transverse overlap area of the colliding nuclei

$\tau$  - is the formation time

[2] - J.T. Mitchell (for the PHENIX Collaboration), Transverse Energy Measurements from the Beam Energy Scan in PHENIX, Nuclear Physics A 00 (2022)

6.

# Calculating the average $p_{\perp}$

The Levy function:

$$f_L(p_{\perp}) = \frac{norm \cdot (n-1)(n-2)}{nT(nT + m_0(n-2))} \cdot p_{\perp} \left( 1 + \frac{\sqrt{m_0^2 + p_{\perp}^2} - m_0}{nT} \right)^{-n}$$

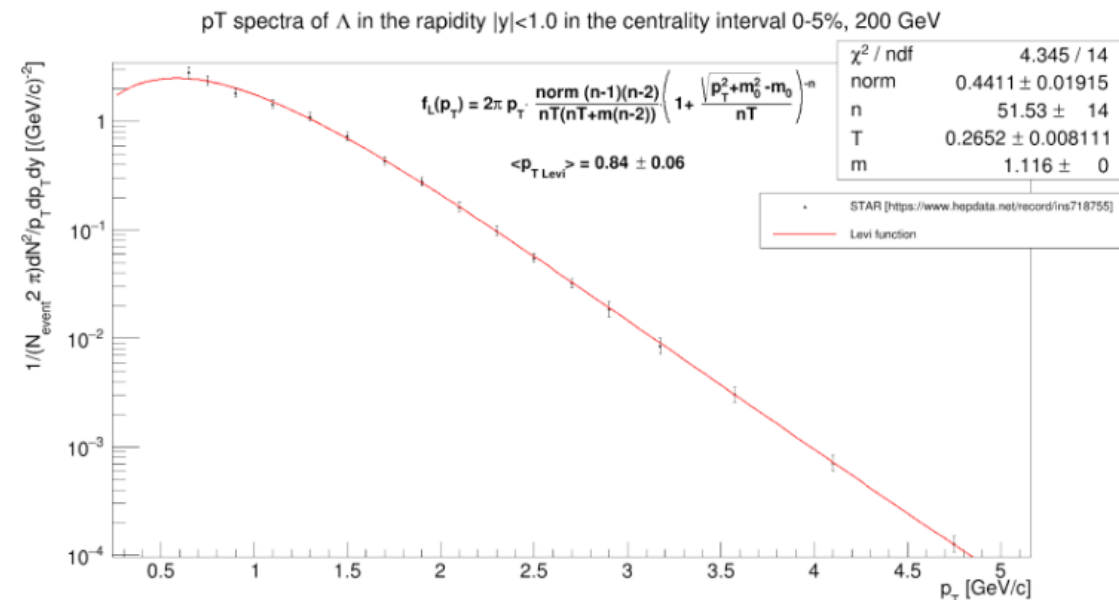
$n$ ,  $T$ , and  $m_0$  - fitting parameters

Average transverse momentum:

$$\langle p_{\perp} \rangle = \frac{\int_a^b f(p_{\perp}) p_{\perp} dp_{\perp}}{\int_a^b f(p_{\perp}) dp_{\perp}}$$

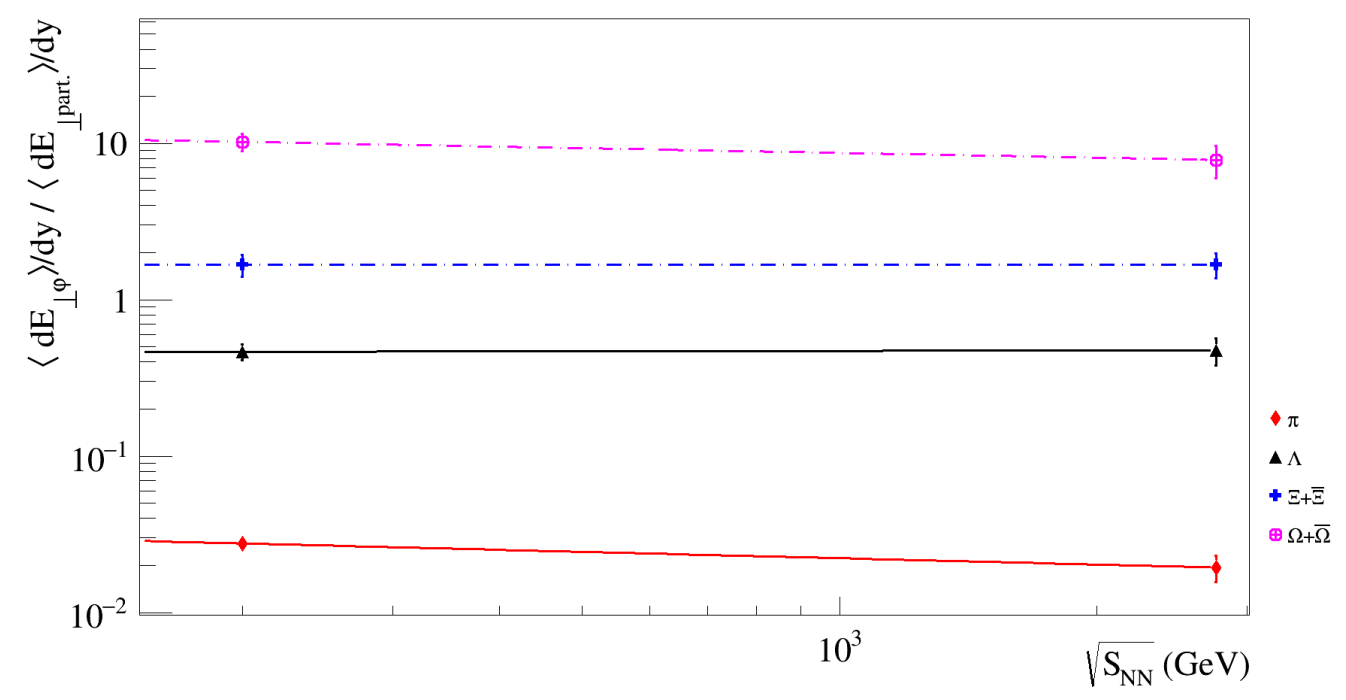
$a$  and  $b$  are extended intervals over the entire range of  $p_{\perp}$

To estimate systematic errors we used the blast-wave function.



$$\langle dE_{\perp}/dy \rangle_{\phi} / \langle dE_{\perp}/dy \rangle_{\text{particles}}$$

$$\langle dE_{\perp}/dy \rangle_{\phi} / \langle dE_{\perp}/dy \rangle_{\text{particles}} = Q * (\sqrt{S_{NN}})^n$$



*A different energies, the fractions of energy expended on the production of various particles with strangeness remain constant.*

Предлагаю вставить картинку только с фи-мезон и омега

# Theoretical models

## Schwinger-like formula

$$|\mu^2| \sim \exp\left(\frac{-\pi m^2}{k}\right)$$

[10]

$k$  is the string tension

where  $M$  gives the matrix-element for the production of a quark-antiquark pair.

$$k \approx 1 \text{ GeV/fm}$$

[10] S.A. Bass et al. “Microscopic Models for Ultrarelativistic Heavy Ion Collisions”, Nuclear Theory (nucl-th); High Energy Physics - Phenomenology

## Thermal model

$$\left(\frac{dE_m}{dy}\right)_{y=0} = \frac{g_m V \lambda_m}{(2\pi)^2} \int \frac{m_T^3 dm_T}{\left[\exp\left(\frac{m_T}{T}\right)\right]}$$

$T$  – is temperature

[11]

$\langle p \rangle$  — mean string tension.

[11] V.V. Vechernin, Physics of Particles and Nuclei, 2023, Vol. 54, No. 3, pp. 528–535, 2023.

# Schwinger-like formula

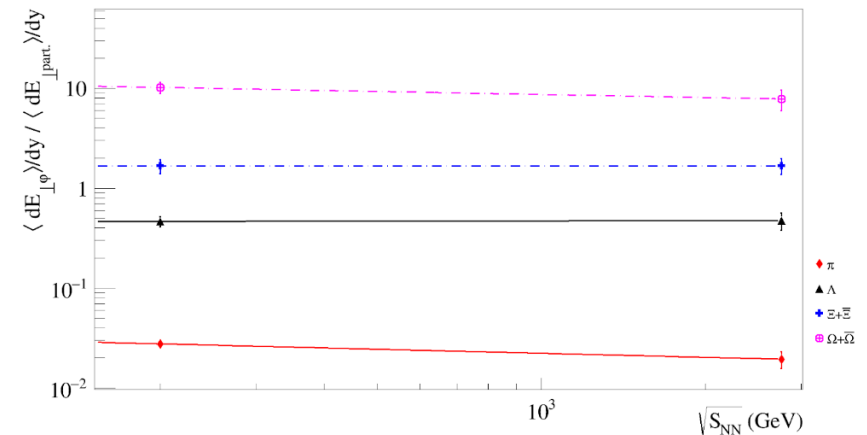
$$\frac{dN}{dy} \sim |\mu|^2 \sim \exp\left(-\frac{\pi \langle m_T \rangle^2}{k}\right)$$

$$\frac{\langle \frac{dE_\varphi}{dy} \rangle}{\langle \frac{dE_\Omega}{dy} \rangle} = \frac{\frac{dN_\varphi}{dy} \langle m_T \rangle_\varphi}{\frac{dN_\Omega}{dy} \langle m_T \rangle_\Omega} = \frac{(2S_\varphi + 1) \langle m_T \rangle_\varphi}{(2S_\Omega + 1) \langle m_T \rangle_\Omega} \exp\left(\frac{\langle m_T \rangle_\Omega^2 - \langle m_T \rangle_\varphi^2}{k} \cdot \pi\right) = 10$$

$$\exp\left(\frac{\langle m_T \rangle_\Omega^2 - \langle m_T \rangle_\varphi^2}{k} \cdot \pi\right) = \frac{10 (2S_\Omega + 1) \langle m_T \rangle_\Omega}{(2S_\varphi + 1) \langle m_T \rangle_\varphi}$$

$$k = \frac{\pi (\langle m_T \rangle_\Omega^2 - \langle m_T \rangle_\varphi^2)}{\ln\left(\frac{10 (2S_\Omega + 1) \langle m_T \rangle_\Omega}{(2S_\varphi + 1) \langle m_T \rangle_\varphi}\right)}$$

$\sqrt{S_{NN}}$	k
39	3,13
200	2,82
2760	4,69





# Thermal model

$$\left(\frac{dE_m}{dy}\right)_{y=0} = \frac{g_m V \lambda_m}{(2\pi)^2} \int \frac{m_T^3 dm_T}{\left[\exp\left(\frac{m_T}{T}\right)\right]}$$

$$\langle E_T \rangle = \left(\frac{\pi}{8} + \frac{1}{4}\right) \left[\langle E \rangle - m_N \langle N_B - N_{\bar{B}} \rangle\right].$$

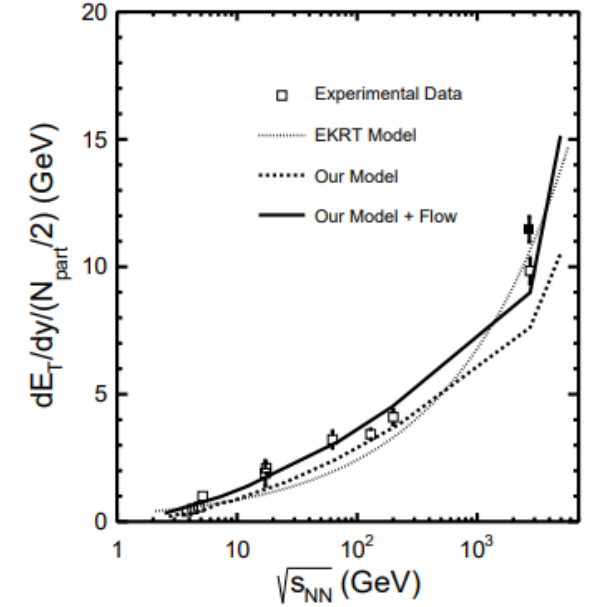
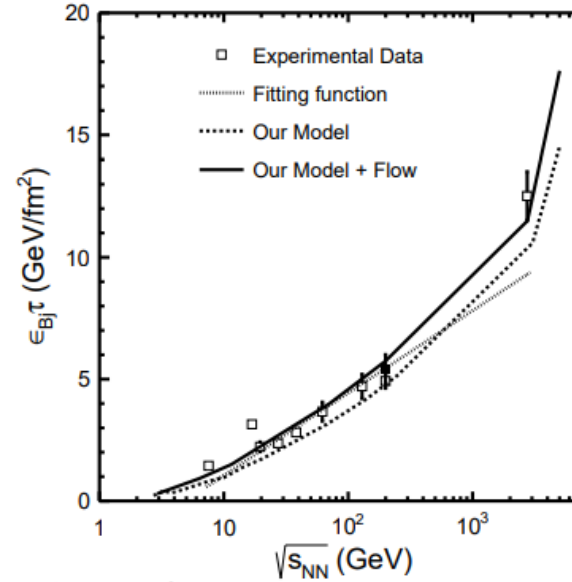


Table 2: Temperature and Baryon Chemical Potential extracted after fitting the particle ratios for various centrality at  $\sqrt{s_{NN}} = 200$  GeV and 2.76 TeV. Most-central, mid-central and peripheral are defined in the text for both energies.

Centrality	$\sqrt{s_{NN}} = 200$ GeV		$\sqrt{s_{NN}} = 2.76$ TeV	
	$T$ (MeV)	$\mu_B$ (MeV)	$T$ (MeV)	$\mu_B$ (MeV)
Most-central	169	23.5	169	1.7
Mid-central	168.5	17	168.5	1.0
Peripheral	168	5.5	168	0.5

[1] - Swatantra Kumar Tiwari , Raghunath Sahoo, “Transverse Energy per Charged Particle in Heavy-Ion Collisions: Role of Collective Flow”, Eur. Phys. J. A (2018) 54: 39

# Что есть на данный момент?

generacy factor:

$$g_i = (2s + 1)(2I + 1).$$

fugacity:

$$\lambda_m = \exp\left(\frac{\mu_m}{T}\right)$$

$$\mu_h = B_h \mu_{Bh} + Q_h \mu_{Qh} + S_h \mu_{Sh},$$

Фактор вырождения  $g$  для частицы  $i$  учитывает все возможные внутренние степени свободы, которые могут влиять на её состояние.

Формула, используемая в модели SHGM:

$$\left(\frac{dE_m}{dy}\right)_{y=0} = \frac{g_m V \lambda_m}{(2\pi)^2} \int \frac{m_T^3 dm_T}{\left[\exp\left(\frac{m_T}{T}\right)\right]}.$$

Преобразованная формула:

$$\left(\frac{dE_m}{dy}\right)_{y=0} = \frac{g_m V \lambda_m}{(2\pi)^2} \cdot \left( e^{-\frac{m_T}{T}} (-m_T^3 T - 3m_T^2 T^2 - 6m_T T^3 - 6T^4) + e^{-\frac{m_T}{T}} \cdot const \right)$$

Для отношения  $\phi$  к  $\Omega$ :

$$\frac{\frac{dE_\phi}{dy}}{\frac{dE_\Omega}{dy}} = \frac{g_\phi \lambda_\phi}{g_\Omega \lambda_\Omega} \cdot \frac{\left( e^{-\frac{m_{T\phi}}{T}} (-m_{T\phi}^3 T - 3m_{T\phi}^2 T^2 - 6m_{T\phi} T^3 - 6T^4) + e^{-\frac{m_{T\phi}}{T}} \cdot const \right)}{\left( e^{-\frac{m_{T\Omega}}{T}} (-m_{T\Omega}^3 T - 3m_{T\Omega}^2 T^2 - 6m_{T\Omega} T^3 - 6T^4) + e^{-\frac{m_{T\Omega}}{T}} \cdot const \right)}$$

# Некоторые преобразования ф-лы

$$\frac{\frac{dE_\phi}{dy}}{\frac{dE_\Omega}{dy}} = \frac{g_\phi \lambda_\phi \cdot \left( e^{\frac{-m_{T\phi}}{T}} (-m_{T\phi}^3 T - 3m_{T\phi}^2 T^2 - 6m_{T\phi} T^3 - 6T^4) + e^{\frac{-m_{T\phi}}{T}} \cdot const \right)}{g_\Omega \lambda_\Omega \cdot \left( e^{\frac{-m_{T\Omega}}{T}} (-m_{T\Omega}^3 T - 3m_{T\Omega}^2 T^2 - 6m_{T\Omega} T^3 - 6T^4) + e^{\frac{-m_{T\Omega}}{T}} \cdot const \right)}$$

Используемые допущения и значения T(статья) и  $\sqrt{s_{NN}}$ :

$$\mu_\phi = \mu_\Omega = \mu_B \rightarrow \lambda_\phi = \lambda_\Omega$$

$$\frac{\frac{dE_\phi}{dy}}{\frac{dE_\Omega}{dy}} = 1$$



$$\frac{g_\Omega \cdot \exp(\mu_\Omega/T)}{g_\phi} = \frac{e^{\frac{-m_{T\phi}}{T}} (-m_{T\phi}^3 T - 3m_{T\phi}^2 T^2 - 6m_{T\phi} T^3 - 6T^4)}{e^{\frac{-m_{T\Omega}}{T}} (-m_{T\Omega}^3 T - 3m_{T\Omega}^2 T^2 - 6m_{T\Omega} T^3 - 6T^4)} = 8,9514$$

$$const = 0$$

$$T = 169 \text{ MeV}$$

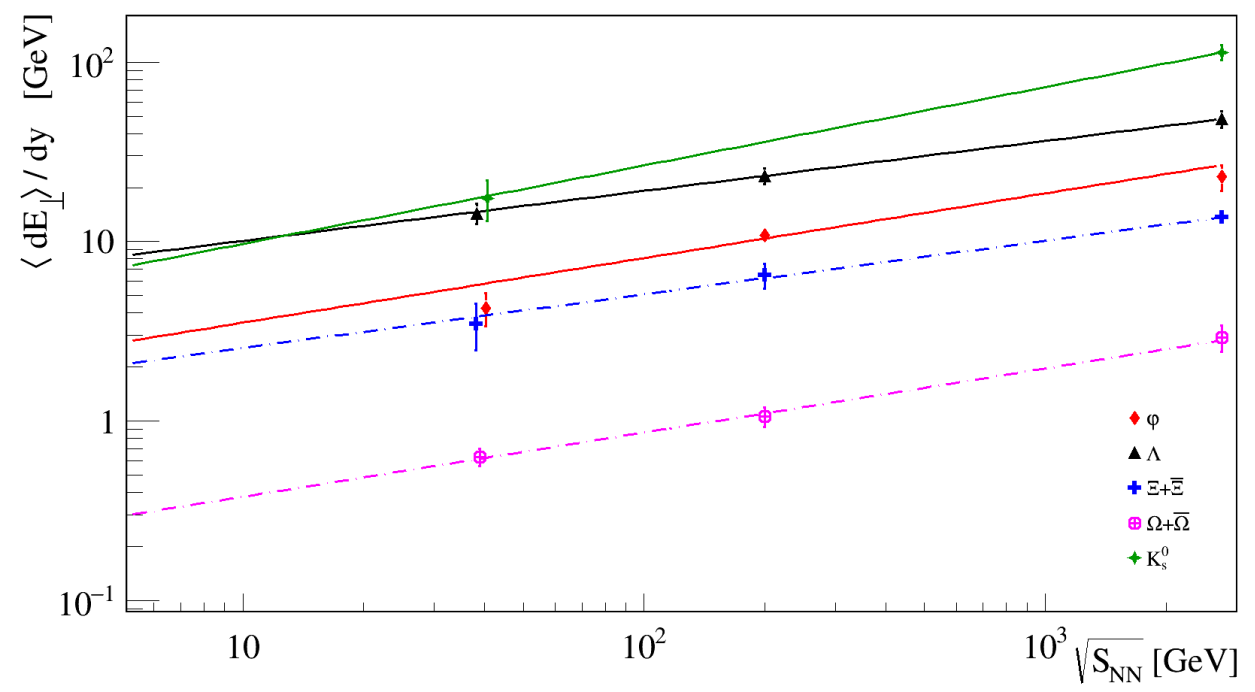
$$\sqrt{s_{NN}} = 200 \text{ GeV}$$

# **Back-up**

**Back-up Back-up Back-up Back-up Back-up Back-up Back-upvvvvvvv Back-up**

8.

# $\langle dE_{\perp}/dy \rangle$ vs. $\sqrt{s_{NN}}$ , very central (0-5 %) $|\eta| < 0.5$



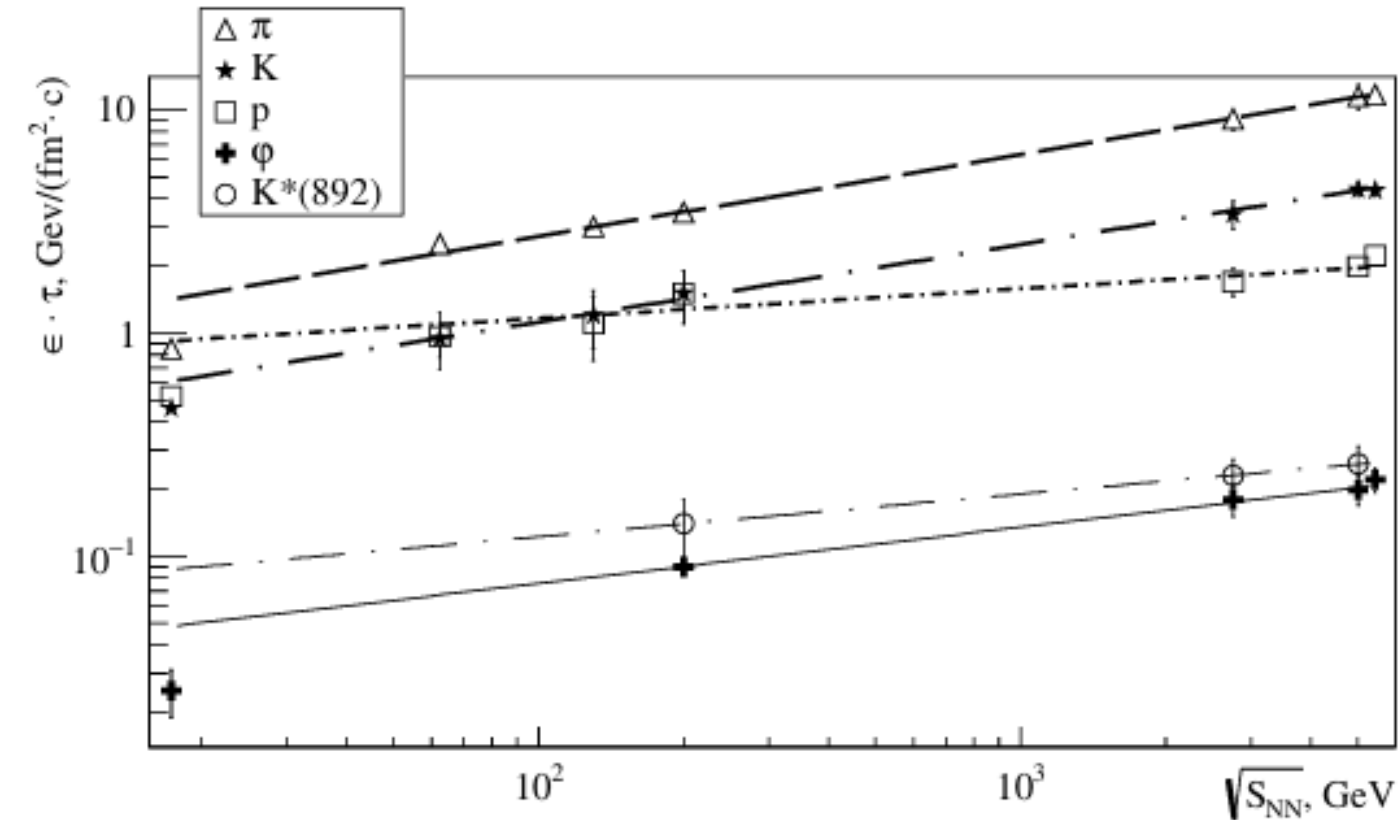
**Approximation function:  $\langle dE_{\perp}/dy \rangle = Q * (\sqrt{s_{NN}})^n$**

Particles	n
$\phi$	$0.36 \pm 0.04$
$\Xi + \bar{\Xi}$	$0.30 \pm 0.05$
$\Lambda$	$0.28 \pm 0.04$
$\Omega + \bar{\Omega}$	$0.39 \pm 0.08$
$K_0^s$	$0.44 \pm 0.06$

The dependence of  $\langle dE_{\perp}/dy \rangle$  on  $\sqrt{s_{NN}}$  for the most central collisions (0-5%) at energies of 39, 200, and 2760 GeV.

*Similar functional dependencies for different particles.*

# $\phi$ -meson and particles with u- and d- quarks [9]

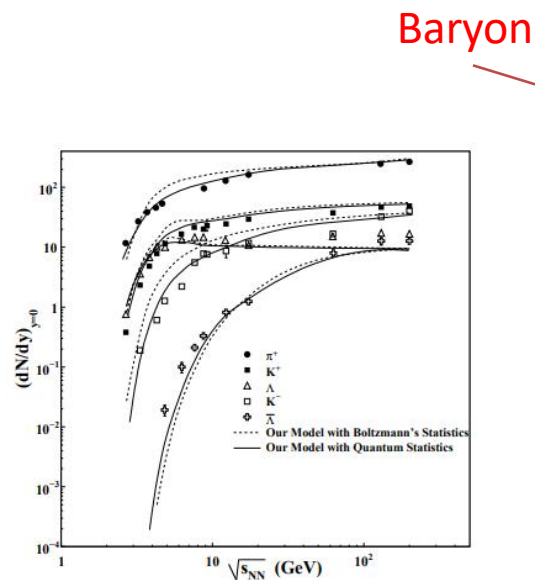


Approximation function:  $\langle dE_{\perp}/dy \rangle = Q$   
 $\cdot (vS_{NN})^n$

Particles	n/2
$\phi$	$0.18 \pm 0.02$
$\pi$	$0.184 \pm 0.05$

The slope of the function for  $\pi$  mesons is similar from that of strange hadrons.

# Thermal model: $dN/dy$



$$\left(\frac{dN_i}{dy}\right)_{th} = \frac{g_i V \lambda_i}{(2\pi^2)} \left[ \left( (1-R) - \lambda_i \frac{\partial R}{\partial \lambda_i} \right) \int_0^\infty \frac{m_T^2 \cosh y \, dm_T}{\left[ \exp\left(\frac{m_T \cosh y}{T}\right) + \lambda_i \right]} - \lambda_i (1-R) \int_0^\infty \frac{m_T^2 \cosh y \, dm_T}{\left[ \exp\left(\frac{m_T \cosh y}{T}\right) + \lambda_i \right]^2} \right].$$

$$\left(\frac{dN_m}{dy}\right)_{th} = \frac{g_m V \lambda_m}{(2\pi^2)} \int_0^\infty \frac{m_T^2 \cosh y \, dm_T}{\left[ \exp\left(\frac{m_T \cosh y}{T}\right) - \lambda_m \right]}.$$

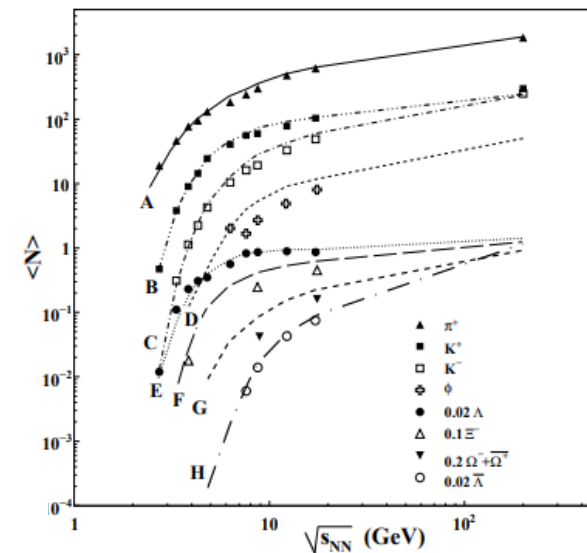


FIG. 3: Variation of rapidity distributions of various hadrons with respect to  $\sqrt{s_{NN}}$  at midrapidity.

Lines show our model calculation. Symbols are the experimental data [9, 49, 50].

$$\frac{dN_i}{dy} = \int_{-\eta_{max.}}^{\eta_{max.}} \left(\frac{dN_i}{dy}\right)_{th} (y - \eta) \, d\eta,$$

The resulting rapidity spectra of  $i$ th hadron, after incorporation of the flow velocity in the longitudinal direction is [11, 12]:

FIG. 2: Variations of total multiplicities of  $\pi^+$ ,  $K^+$ ,  $K^-$ ,  $\phi$ ,  $\Lambda$ ,  $\Xi^-$ ,  $(\Omega^- + \Omega^+)$ , and  $\bar{\Lambda}$  with respect to center-of-mass energy predicted by our model. Experimental data measured in central  $Au - Au/Pb - Pb$  collisions [31–47] have also been shown for comparison. In this figure, A, B, C, D, E, F, G, and H represent the multiplicities of  $\pi^+$ ,  $K^+$ ,  $K^-$ ,  $\phi$ ,  $\Lambda$ ,  $\Xi^-$ ,  $(\Omega^- + \Omega^+)$ , and  $\bar{\Lambda}$ , respectively.

# Sources of experimental data used

GeV	Particles	GeV	GeV	GeV	GeV	GeV
		GeV	GeV	GeV	GeV	GeV
62.4	Pions	0.4 [1]	23.7 ± 1.1 [1]	133 ± 17 [1]	188.7 ± 14.3 [1]	23 ± 0.16 [1]
	Kaons	0.6-0.65 ± 0.05 [1]	32.4 ± 2.3 [1]	37.6 ± 2.7 [1]		111.5 ± 16.2 [1]
	Protons	0.95 ± 0.1 [1]	13.6 ± 1.7 [1]	29.0 ± 3.8 [1]		113.8 ± 26.8 [1]
130	Pions	0.4 [1]	280 ± 25 [1]	278 ± 25 [1]		355 ± 21 [1]
	Kaons	0.65-0.7 ± 0.05 [1]	42.7 ± 6.2 [1]	46.3 ± 6.5 [1]		145.8 ± 35.3 [1]
	Protons	1 ± 0.1 [1]	20.0 ± 3.4 [1]	28.2 ± 4.4 [1]		132.2 ± 35.4 [1]
200	Pions	0.4 [1]	327 ± 25 [1]	322 ± 25 [1]		392 ± 24 [1]
	Kaons	0.7-0.8 ± 0.05 [1]	49.5 ± 6.2 [1]	51.3 ± 6.5 [1]		181 ± 40.8 [1]
		1.08 [2]	10.482 ± 4 [2]			14.74 ± 3 [1]
	Protons	1.1 ± 0.1 [1]	26.7 ± 3.4 [1]	34.7 ± 4.4 [1]		177.6 ± 41.3 [1]
		0.970 ± 0.2 [3]	7.70 ± 3 [3]			10.840 ± 5 [1]
		0.96 ± 0.1 2 [8]	0.9 ± 0.1 [8]	1.83 ± 0.2 [6]	2.17 ± 0.19 [6]	2.98 ± 0.5 3.5 ± 0.5
		0.84 ± 0.0 6 [8]	0.85 ± 0.0 [8]	16.7 ± 1.1 [6]	12.7 ± 0.9 [6]	23.3 ± 2.4 17.8 ± 1.9
		1.095 ± 0.111 [8]		0.53 ± 0.04 [6]		1.06 ± 0.13
2760	Pions	0.517 ± 0.019 [4]	0.520 ± 0.018 [4]	733 ± 54 [4]	732 ± 52 [11]	1179.96 ± 110.7
	Kaons	0.876 ± 0.026 [4]	0.867 ± 0.027 [4]	109 ± 9 [4]	109 ± 9 [11]	436.7 ± 56.5
		1.310 ± 0.06 [5]	19.562 ± 6 [5]			30.95 ± 0.9 30.95 ± 0.9 30.95 ± 0.9 30.95 ± 0.9
	Protons	1.333 ± 0.033 [4]	1.353 ± 0.034 [4]	34 ± 3 [4]	33 ± 3 [4]	219.5 ± 27.1
		1.310 ± 0.07 [6]	13.81 ± 8 [6]			22.93 ± 8 0.180 ± 0.3 22.93 ± 8
		1.57 ± 0.0 4 [10]	1.56 ± 0.0 4 [10]	6.67 ± 0.47 [3]		13.66 ± 1.02
		1.48 ± 0.04 [4]		26 ± 3 [4]		48.19 ± 5
		1.8 ± 0.1 [10]		1.19 ± 0.19 [3]		2.92 ± 0.49
5020	Pions	0.5682 [7]	1699.80 [7]			1491.8 ± 167.2
	Kaons	0.9177 [7]	273.41 [7]			569.8 ± 34.8
		1.460 ± 0.07 [8]	19.72 ± 8 [8]			37.54 ± 6.5 0.290 ± 0.5
	Protons	1.4482 [7]	74.56 [7]			257 ± 18 1.99 ± 0.16
	1.4 ± 0.02 [8]	14.9 ± 1.2 [8]			29.83 ± 9 0.23	
5044	Pions	0.530 ± 0.02 [9]	1002.6757 ± 2 [9]			826.6551 ± 5
	Kaons	0.90 ± 0.03 [9]	149.3714 ± 0.7 [9]			308
	Protons	1.40 ± 0.02 [9]	46.214 ± 0.7 [9]			156.9
		1.330 ± 0.03 [9]	9.271 [9]			31.51 ± 9 0.430 ± 0.3

Я таки не знаю как её вставлять



# Индусы, статья 3.0

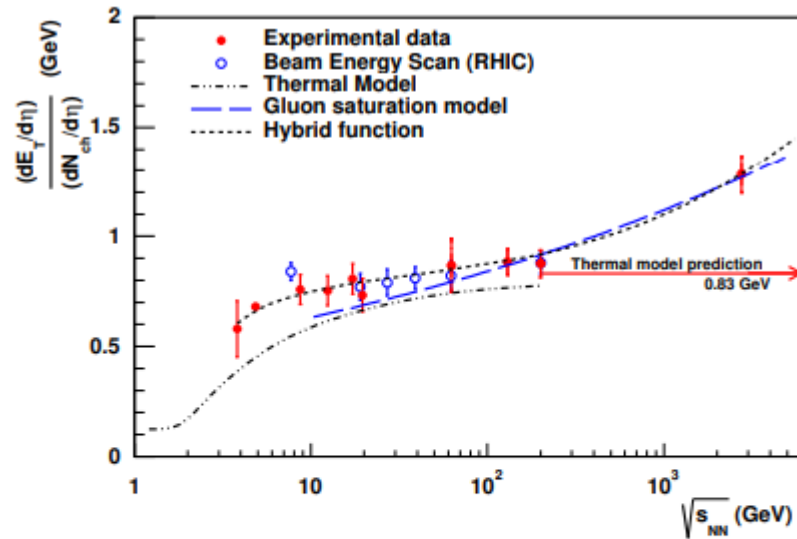


FIG. 42: The ratio of  $\frac{dE_T}{d\eta}$  and  $\frac{dN_{ch}}{d\eta}$  at midrapidity, as a function of center of mass energy. Experimental data are compared to the predictions from thermal model, gluon saturation model and the estimations obtained in the framework of the hybrid model fitting to transverse energy and charged particle data.

# Нормальный слайд давай сделаем

## Тезисы введения:

- 1) Есть модель SHGM -Statistical Hadron Gas Model
- 2) Для энергий RHIC она хорошо описывает некоторые эффекты???
- 3) Для энергий LHC наблюдаются отклонения, например отклонения зависимости  $dN_{ch}/d\eta$  от энергии от логарифмической
- 4) В статье используется модель SHGM с включенным потоком, для объяснения этих эффектов.

# 모든 종류의 헛소리

Freeze-Out Parameters in Heavy-Ion Collisions at AGS, SPS, RHIC, and LHC Energies

????

$$\langle E_T \rangle = \left( \frac{\pi}{8} + \frac{1}{4} \right) \left[ \langle E \rangle - m_N \langle N_B - N_{\bar{B}} \rangle \right].$$

"Relativistic Heavy-Ion Physics" by Reinhard Stock:

$$\mu_h = B_h \mu_{Bh} + Q_h \mu_{Qh} + S_h \mu_{Sh}$$

$$m_T = \sqrt{m^2 + p_T^2}$$

$$\left( \frac{dE_m}{dy} \right)_{y=0} = \frac{g_m V \lambda_m}{(2\pi)^2} \int \frac{m_T^3 dm_T}{\left[ \exp \left( \frac{m_T}{T} \right) \right]}$$

$$g_i = (2s_i + 1)(2I_i + 1),$$

Centrality	$\sqrt{s_{NN}} = 200$ GeV		$\sqrt{s_{NN}} = 2.76$ TeV	
	$T$ (MeV)	$\mu_B$ (MeV)	$T$ (MeV)	$\mu_B$ (MeV)
Most-central	169	23.5	169	1.7
Mid-central	168.5	17	168.5	1.0
Peripheral	168	5.5	168	0.5

$$z_i = e^{\mu_i/T}$$

$$-x^3 \cdot T \cdot e^{-\frac{x}{T}} - T^2 \cdot 3 \cdot x^2 \cdot e^{-\frac{x}{T}} - T^3 \cdot 6 \cdot x \cdot e^{-\frac{x}{T}} - T^4 \cdot 6 \cdot e^{-\frac{x}{T}} + Const$$