



C, P, T conjugation formulas for spinor field operators

Vadim Monakhov

Associate Professor, v.v.monahov@spbu.ru

Saint-Petersburg State University

LXXV International conference Nucleus-2025.

Saint Petersburg, Russia. 1-6 July, 2025.

CPT symmetry

CPT symmetry plays a key role in modern quantum field theory (QFT). It guarantees for particles and antiparticles equality of masses, decay widths, equality in absolute value and opposite sign of all types of charges.

CPT:

- is **anti-unitary** operator,
- preserves the Dirac equation $\gamma^\mu(i\partial_\mu - qA_\mu(x))\psi(x) = m\psi(x)$,
- particle \rightarrow antiparticle, charge $q \rightarrow -q$, potential $A_\mu \rightarrow -A_\mu$,
- energy $p_0 \rightarrow p_0$,
- momentum $p_k \rightarrow p_k$ (i.e. $\mathbf{p} \rightarrow \mathbf{p}$),
- spin projection $s_3 \rightarrow -s_3$,
- **chirality** $\mathbf{L} \rightarrow \mathbf{R}$, $\mathbf{R} \rightarrow \mathbf{L}$, where $\psi_L = \frac{1 - \gamma^5}{2} \psi$, $\psi_R = \frac{1 + \gamma^5}{2} \psi$

(L - left-chiral, R - right-chiral. Chirality and helicity should not be confused!)

General properties of C, P, T conjugations (inversions)

Operators C, P, T

- preserve the Dirac equation $\gamma^\mu(i\partial_\mu - qA_\mu) \psi = m\psi$,
- preserve energy $E = p_0 \rightarrow p_0$.

C: unitary operator, which transforms

- particle \rightarrow antiparticle, charge $q \rightarrow -q$; potential $A_\mu \rightarrow -A_\mu$,
- momentum $\mathbf{p} \rightarrow \mathbf{p}$,
- spin projection $s_3 \rightarrow s_3$,
- **preserves chirality** $L \rightarrow L, R \rightarrow R$

P: unitary operator, which transforms

- particle \rightarrow particle; charge $q \rightarrow q$; potential $A_0 \rightarrow A_0, \mathbf{A} \rightarrow -\mathbf{A}$,
- momentum $\mathbf{p} \rightarrow -\mathbf{p}$,
- spin projection $s_3 \rightarrow s_3$,
- **changes chirality** $L \rightarrow R, R \rightarrow L$

T: anti-unitary operator, which transforms

- particle \rightarrow particle; charge $q \rightarrow q$; potential $A_0 \rightarrow A_0, \mathbf{A} \rightarrow -\mathbf{A}$,
- momentum $\mathbf{p} \rightarrow -\mathbf{p}$,
- spin projection $s_3 \rightarrow -s_3$,
- **preserves chirality** $L \rightarrow L, R \rightarrow R$

Erroneous formulas of CPT conjugation

Standard formulas of CPT conjugation are [1], [2], [5]

$$\text{CPT } \psi(t,x) = \eta \gamma^5 \psi(-t,-x), \quad (1)$$

or [2], [3], [4], [6], [7]

$$\text{CPT } \psi(t,x) = \eta \gamma^5 \psi^{+Tmatr}(-t,-x), \quad (2)$$

or [7], [8]

$$\text{CPT } \psi(t,x) = \eta \gamma^5 \psi^*(-t,-x) = \eta \gamma^5 \psi^{+T}(-t,-x), \quad (3)$$

where $(.)^{Tmatr}$ is matrix transposition, and $(.)^T = (.)^{+*} = (.)^+ (.)^*$

- [1]. W. Pauli. Exclusion principle, Lorentz group and reflexion of space-time and charge. In: *Niels Bohr and the Development of Physics*. Pergamon Press, 30-51 (1955).
- [2]. G. Lüders. Proof of the TCP Theorem. *Annals of Physics*, 2(1), 1-15 (1957).
- [3]. G. Feinberg, S. Weinberg. On the Phase Factors in Inversions. *Il Nuovo Cimento*, 14, 571-592 (1959).
- [4]. G. Grawert, G. Lüders, H. Rollnik. The TCP theorem and its applications. *Fortschritte der Physik*, 7(6), 291-328 (1959).
- [5]. J. Bjorken, S. Drell, Relativistic Quantum Mechanics, McGraw-Hill College: N.Y. (1964).
- [6]. J. Bjorken, S. Drell, Relativistic Quantum Fields, McGraw-Hill College: N.Y. (1965).
- [7]. M. Peskin, D. Schroeder. An Introduction to Quantum Field Theory, Addison-Wesley Publishing Company: USA (1995).
- [8]. S. Weinberg. The Quantum Theory of Fields - Foundations. Vol. 1. Cambridge University Press: Cambridge, 1995.

Why these formulas are erroneous?

Let according to (1)

$$\text{CPT } \psi(t, x) = \eta \gamma^5 \psi(-t, -x) ,$$

$$\Rightarrow \text{CPT} = \eta \gamma^5 R(-t, -x) , \text{ where } R(-t, -x) \psi(t, x) = \psi(-t, -x) .$$

$R(-t, -x) = R(-t) R(-x)$ acts only on the t and x components, not on the matrices or creation and annihilation operators - it is very important!

$$\Rightarrow \text{CPT } \gamma^5 = \gamma^5 \text{CPT}$$

$$\Rightarrow \text{CPT } \psi_L(t, x) = \eta \gamma^5 \psi_L(-t, -x) = (\eta \gamma^5 \psi(-t, -x))_L$$

This CPT operator transforms left-chiral spinor to left-chiral antispinor!

The same is for formulas (2) and (3) of the CPT.

According to the Standard Model there should be:

$$\text{CPT } \psi_L = \psi_R^a , \text{ where } \psi^a - \text{field operator of the antiparticle.}$$

\Rightarrow formulas (1), (2), (3) are erroneous!

Where is the error in the formulas for C, P, T?

1. Space inversion operator is unitary and changes chirality $L \rightarrow R, R \rightarrow L$:

$$\mathbf{P} = \eta_P \gamma^0 R (-\mathbf{x}) \quad (4)$$

2. Time inversion operator is anti-unitary and preserves chirality $L \rightarrow L, R \rightarrow R$

$$\mathbf{T} = \eta_T \gamma^1 \gamma^3 R (-t) (.)^* \quad (5)$$

3. There are **four charge conjugation operators** (we use Dirac or Weil representation of gamma matrices): **non-QFT** operators C_{M-K} and C_{Pauli} ; **QFT** operators C_{Schw} and C_{QFT} .

- $C_{M-K} = \eta_C \gamma^2 (.)^*$ – Majorana-Kramers; **changes chirality, anti-unitary \Rightarrow error.**
- C_{Pauli} : C_{M-K} and «hole» \rightarrow antiparticle – **changes chirality, anti-unitary \Rightarrow error.**
- $C_{Schw} = \eta_C \gamma^2 (.)^{+T_{matr}}$ – Schwinger charge conjugation is QFT analog of C_{Pauli} ; **changes chirality, anti-unitary \Rightarrow error;**
- C_{QFT} : $b_s(p) \rightarrow \eta_C d_s(p), d_s^+(p) \rightarrow \eta_C b_s^+(p)$; **preserves chirality, unitary**
 $\Rightarrow C_{QFT}$ is the operator C of the CPT theorem!

The origin of the error: $C_{QFT} \psi = C_{Schw} \psi$ for non-chiral spinors

\Rightarrow erroneous conclusion that “ $C_{QFT} \equiv C_{Schw}$ ”. However, $C_{QFT} \psi_L \neq C_{Schw} \psi_L$.

A well-known implementation of C_{QFT}

J. Bjorken, S. Drell [6] :

$$C_{QFT}\psi(x) = \frac{1}{(2\pi)^{3/2}} \int d^3p \sqrt{\frac{m}{E}} \left(C b_s(p) C^{-1} u_s(p) e^{-ip_\mu x^\mu} + C d_s^+(p) C^{-1} v_s(p) e^{ip_\mu x^\mu} \right)$$

$$C = e^{iA_1\varphi} e^{-iA_2\varphi} = e^{-iA_2\varphi} e^{iA_1\varphi}; \varphi = \frac{\pi}{2}$$

$$A_1 = A_1^+ = \int d^3p (b_s^+(p)b_s(p) + d_s^+(p)d_s(p));$$

$$A_2 = A_2^+ = \int d^3p (b_s^+(p)d_s(p) + d_s^+(p)b_s(p));$$

$$C_{QFT} = \int d^3p [b_s^+(p)b_s(p) + d_s^+(p)d_s(p), \bullet] \int d^3p [b_s^+(p)d_s(p) + d_s^+(p)b_s(p), \bullet]$$

$$C_{QFT}b_s(p) = Cb_s(p)C^{-1} = d_s(p); \quad C_{QFT}b_s^+(p) = d_s^+(p);$$

$$C_{QFT}d_s^+(p) = Cd_s^+(p)C^{-1} = b_s^+(p); \quad C_{QFT}d_s(p) = b_s(p);$$

Errors in the formulas for P and T

- Is operator $\mathbf{C}_{\text{QFT}} \mathbf{PT} = C_{\text{QFT}} \eta_P \eta_T R(-t, -x) i \gamma^2 \gamma^5 (.)^*$? It is anty-unitary, changes chirality $L \rightarrow R, R \rightarrow L \Rightarrow$ is all correct? – **No**. Why?
- Previous authors (J. Bjorken, S. Drell; M. Peskin, D. Schroeder; S. Weinberg and so on) have suggested that the transformation of the spinor field operator can be described by
 - either matrix operators $S(p)$,
 - or transformations of operators $b_s(p), d_s^+(p)$,and that these methods are equivalent.
- Let us prove that this assumption is wrong.

Proper Lorentz transformations

To obtain $\psi_s(p')$ from $\psi_s(p)$, it is necessary to apply active Lorentz transformation.

$$\psi_s(p) = \frac{1}{(2\pi)^{3/2}} \sqrt{\frac{m}{E(p)}} \left(b_{<s>}(p) u_{<s>}(p) e^{-ip_\mu x^\mu} + d_{<s>}^+(p) v_{<s>}(p) e^{ip_\mu x^\mu} \right). \quad (6)$$

$$\psi_s(p') = \frac{1}{(2\pi)^{3/2}} \sqrt{\frac{m}{E(p')}} \left(b_{<s>}(p') u_{<s>}(p') e^{-ip'_\mu x^\mu} + d_{<s>}^+(p') v_{<s>}(p') e^{ip'_\mu x^\mu} \right). \quad (7)$$

\Rightarrow

$$b_s(p) \rightarrow b_s(p'); \quad u_s(p) \rightarrow u_s(p'); \quad e^{-ip_\mu x^\mu} \rightarrow e^{-ip'_\mu x^\mu}, \quad i\partial_\mu e^{-ip'_\mu x^\mu} = p'_\mu,$$

$$d_s^+(p) \rightarrow d_s^+(p'); \quad v_s(p) \rightarrow v_s(p'); \quad e^{ip_\mu x^\mu} \rightarrow e^{ip'_\mu x^\mu}, \quad i\partial_\mu e^{ip'_\mu x^\mu} = -p'_\mu.$$

After the Lorentz transformation $p_\mu x^\mu \rightarrow p'_\mu x'^\mu$. However, we observe the field operator (9) at the same point $x' = x$ and not in the point $x' = \Lambda x$.

That is why $p'_\mu x'^\mu = p'_\mu x^\mu$, and the field operator $\psi(x)$ transforms **into $\psi'(x)$, and not into $\psi'(\Lambda x)$.**

Error in the formula for P

$\psi(x)$ symmetrically contains terms with p and $-p$, s and $-s$. Therefore, it is necessary to investigate transformations of $\psi_s(p)$, and **not** $\psi(x)$.

$$\psi(x) = \frac{1}{(2\pi)^{3/2}} \int d^3 p \sqrt{\frac{m}{E}} \left(b_s(p) u_s(p) e^{-ip_\mu x^\mu} + d_s^+(p) v_s(p) e^{ip_\mu x^\mu} \right) = \int d^3 p \sum_s \psi_s(p), \quad (8)$$

$$\begin{aligned} P \psi_s(p) &= \frac{\eta_P}{(2\pi)^{3/2}} \sqrt{\frac{m}{E(p)}} \left(b_{<s>}(p) u_{<s>}(-p) e^{-ip_\mu \tilde{x}^\mu} - d_{<s>}^+(p) v_{<s>}(-p) e^{ip_\mu \tilde{x}^\mu} \right) = \\ &= \frac{\eta_P}{(2\pi)^{3/2}} \sqrt{\frac{m}{E(\tilde{p})}} \left(b_{<s>}(p) u_{<s>}(\tilde{p}) e^{-i\tilde{p}_\mu x^\mu} - d_{<s>}^+(p) v_{<s>}(\tilde{p}) e^{i\tilde{p}_\mu x^\mu} \right) \neq \eta_P \psi_s(-p), \end{aligned} \quad (9)$$

where $\tilde{p}^0 = p^0$; $\tilde{p}^k = -p^k$, $k = 1, 2, 3$.

$\Rightarrow P \psi_s(p)$ is **invalid field operator** with $b_s(p)$ and $-d_s^+(p)$ instead of $b_s(-p)$ and $-d_s^+(-p)$

$\Rightarrow P$ must transform $b_s(p)$ to $b_s(-p)$ and $d_s^+(p)$ to $-d_s^+(-p)$!

Space inversion operator P

Thus, the action of the spatial inversion operator on the spinor field operator leads to

$$\psi'_s(p) = \frac{\eta_P}{(2\pi)^{3/2}} \sqrt{\frac{m}{E(\tilde{p})}} \left(b_{<s>}(\tilde{p}) u_{<s>}(\tilde{p}) e^{-i\tilde{p}_\mu x^\mu} + d_{<s>}^+(\tilde{p}) v_{<s>}(\tilde{p}) e^{i\tilde{p}_\mu x^\mu} \right),$$

$$\psi'(x) = \eta_P \int d^3\tilde{p} \sum_s \psi_s(\tilde{p}) = \eta_P \int d^3p \sum_s \psi_s(p) =$$

$$= \frac{\eta_P}{(2\pi)^{3/2}} \int d^3p \sqrt{\frac{m}{E}} \left(b_s(p) u_s(p) e^{-ip_\mu x^\mu} + d_s^+(p) v_s(p) e^{ip_\mu x^\mu} \right) = \eta_P \psi(x).$$

$$\Rightarrow \quad P \psi_s(p) = \eta_P \psi_s(-p),$$

$$P \psi(x) = \eta_P \psi(x),$$

$$P \psi_L(x) = \eta_P \psi_R(x),$$

$$P = \eta_P \gamma^0 R(\Lambda),$$

where $\eta_P = \pm 1$ and $R(\Lambda) b_s(p) = b_s(-p)$, $R(\Lambda) d_s^+(p) = -d_s^+(-p)$.

Formulas for proper Lorentz transformations and space-time inversions are similar

- $\mathbf{L} = \eta_{\mathbf{L}} \mathbf{S} R(\Lambda) ; \quad S = \eta_{\mathbf{LS}} \exp(\gamma^0 \gamma(p) \omega(p))$
 $R(\Lambda) b_s(p) = \eta_{\mathbf{Lb}} b_s(\Lambda^{-1} p); \quad R(\Lambda) d_s^+(p) = \eta_{\mathbf{Lb}} d_s^+(\Lambda^{-1} p); \quad \eta_{\mathbf{L}} = \eta_{\mathbf{Lb}} \eta_{\mathbf{LS}} = \eta_{\mathbf{Lb}} = \eta_{\mathbf{LS}} = 1$
- $\mathbf{P} = \eta_{\mathbf{P}} \mathbf{S} R(\Lambda); \quad S = \eta_{\mathbf{PS}} \gamma^0$
 $R(\Lambda) b_s(p) = \eta_{\mathbf{Pb}} b_s(\Lambda^{-1} p) = \eta_{\mathbf{Pb}} b_s(-p); \quad R(\Lambda) d_s^+(p) = -\eta_{\mathbf{Pb}} d_s^+(-p); \quad \eta_{\mathbf{P}} = \eta_{\mathbf{Pb}} \eta_{\mathbf{PS}} = \pm 1.$
 $\mathbf{P} \psi_s(p) = \eta_{\mathbf{P}} \psi_s(-p)$
 $\mathbf{P} \psi(x) = \eta_{\mathbf{P}} \psi(x); \quad \mathbf{P}^2 = 1$
- $\mathbf{T} = \eta_{\mathbf{T}} \mathbf{S} (.)^* R(\Lambda); \quad S = \eta_{\mathbf{TS}} \gamma^1 \gamma^3$
 $R(\Lambda) b_s(p) = \eta_{\mathbf{Tb}} b_{-s}(-p); \quad R(\Lambda) d_s^+(p) = \eta_{\mathbf{Td}} d_{-s}^+(-p); \quad \eta_{\mathbf{T}} = \eta_{\mathbf{Tb}} \eta_{\mathbf{TS}}$
 $\mathbf{T} \psi_s(p) = \eta_{\mathbf{T}} (-s) \psi_{-s}(-p); \quad \mathbf{T}^2 = -1$

CPT transformation

$C_{\text{QFT}} \psi \equiv \psi^a$ – antiparticle field operator.

$$\square \quad \text{CPT} = \eta_{\text{CPT}} R(-t, -x) C_{\text{QFT}} \gamma^5 i \gamma^2 (.)^*$$

$$\square \quad \text{CPT} \psi_s(p) = \eta_{\text{CPT}} (-s) \psi^a_{-s}(p)$$

$$\square \quad \text{CPT} \psi_{sL}(p) = \eta_{\text{CPT}} (-s) \psi^a_{-sR}(p)$$

Conclusions

- “Standard” formulas for the CPT transformation are invalid: they transform a left-chiral spinor into a left-chiral antispinor.
- Correct operator C of the charge conjugation is C_{QFT} such as $C_{\text{QFT}} b_s(p) = d_s(p)$; $C_{\text{QFT}} d_s(p) = b_s(p)$.
- Lorentz transformations acts both on the spinor columns u_s, v_s via the matrix operator S , and on the creation and annihilation operators $b_s(p)$ and $d_s^\dagger(p)$ via the transformation of their argument $p_\mu \rightarrow p'_\mu$ and possible phase factor.
- Base on this we derived the correct formulas for **P**, **T** and **CPT** operators which correspond to the Standard model.