

Threshold amplification of $X(a,b)Y$ reactions: from μCF to EM formfactors of hyperons

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LXXV International conference «NUCLEUS-2025»
St.Peterburg State University, 1-6 July 2025

Problem

Physics highlights at BESIII and STCF

Xiaorong Zhou

University of Science and technology of China

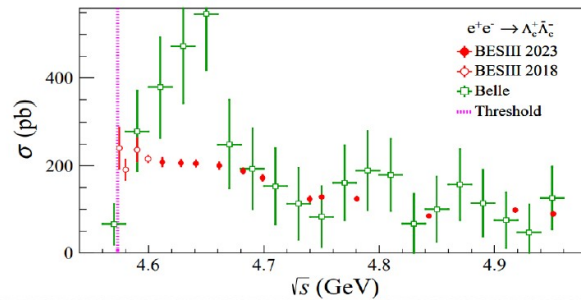
The XXIV Khariton Topical Scientific Readings on Problems of
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25/7/2023, Sarov

Cross section of $e^+e^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-$

M. Ablikim *et al.*

(BESIII Collaboration) PHYSICAL REVIEW LETTERS 131, 191901 (2023)



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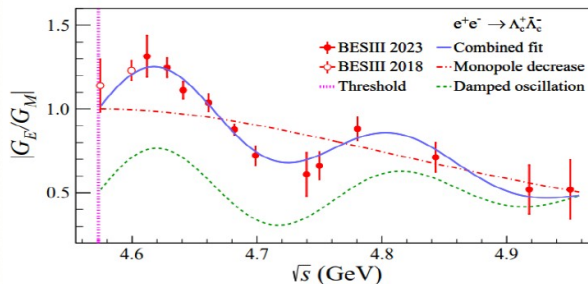
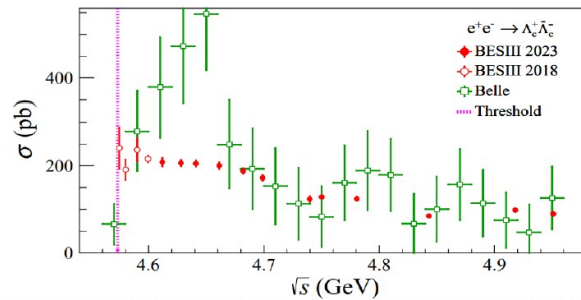
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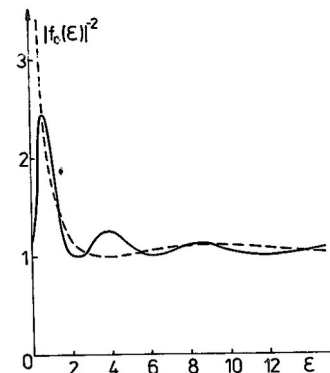
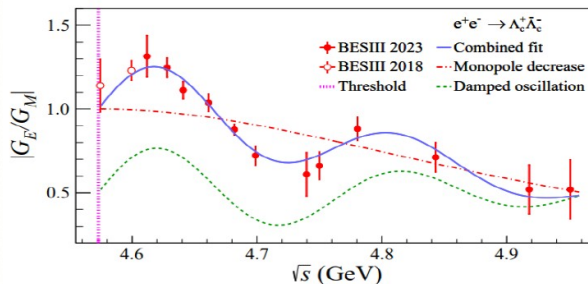
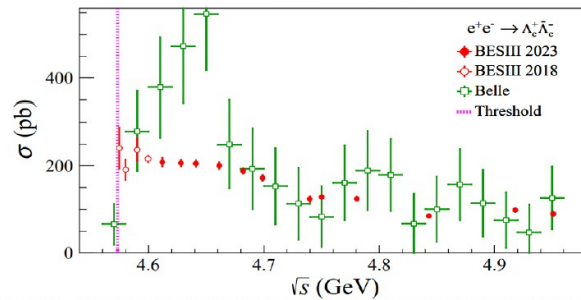
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V.S. Melezhik / Resonance amplification ..
Nuclear Physics A550 (1992) 223–234 ..

Similar problems ?

Nuclear Physics A550 (1992) 223–234
North-Holland

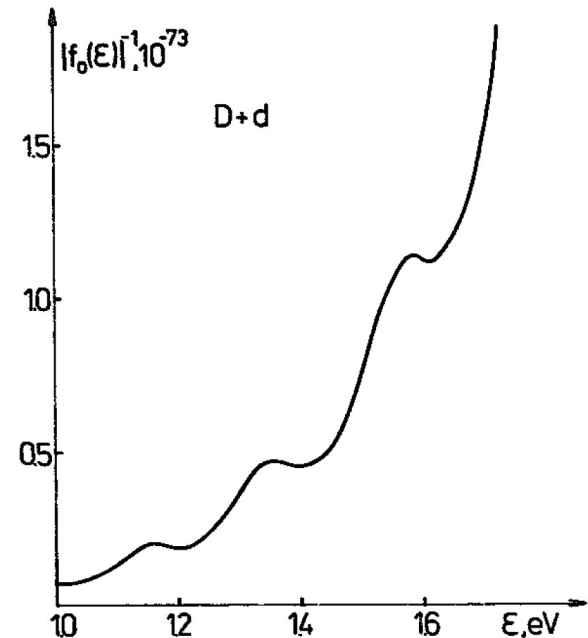
**NUCLEAR
PHYSICS A**

Resonance amplification of the nuclear reaction $X(a, b)Y$ near the $a + X$ channel threshold

V.S. Melezhik

Joint Institute for Nuclear Research, Head P.O. Box 79, Dubna, Moscow Region, Russian Federation

Screening effects in fusion reactions of the type $D(d, p)T$ near the threshold of the channel $d + D$



Similar problems ?

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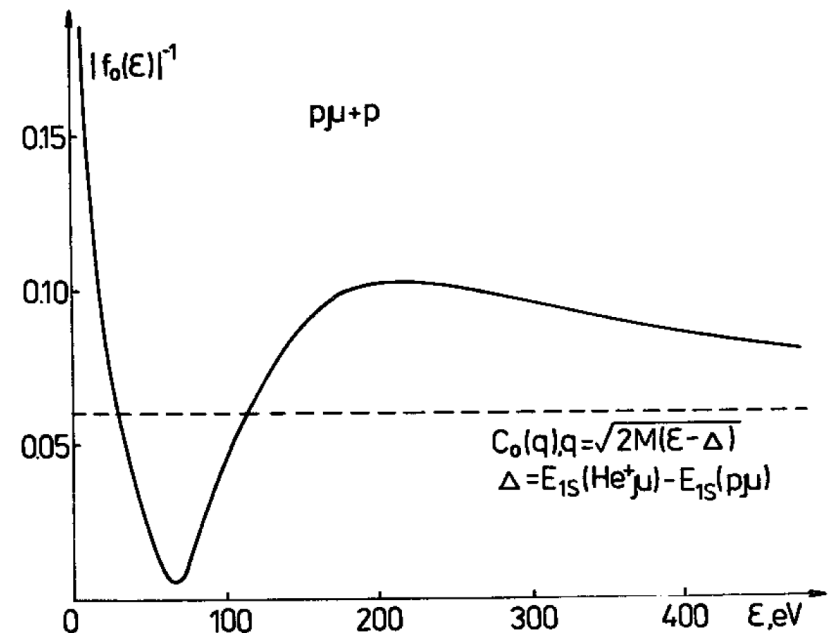
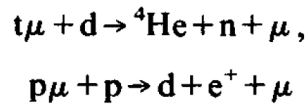
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“In flight” fusion reactions in mesic atomic physics



Simple potential model for $X(a,b)Y$

V.S. Melezhik, Nucl. Phys. A550, 223

(1992)

It is known that the S-wave cross section of the fusion reaction $X(a,b)Y$ is described near the threshold $\varepsilon \rightarrow 0$ of the channel $a + X$ by the formula ¹⁾

$$\sigma^{\text{in}}(\varepsilon) = |\psi_\varepsilon(R=0)|^2 \frac{A}{V}, \quad A = \text{const}, \quad (1)$$

where R is the distance between the fragments a and X , $\psi_\varepsilon(R)$ is the wave function of relative motion of a and X , and $V = \sqrt{2\varepsilon/M}$ is the relative velocity of their motion.

K.R. Leng, Astrophysical formulae, vol.II (Springer, Berlin 1974)

S.Deser, M.L. Goldberger, K. Baumann, W. Thirring, Phys. Rev. 94, 774 (1954)

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If the Coulomb interaction ($1/R$) occurs between the fragments a and X , formula (1) is reduced to the known Gamow formula

$$\sigma_0^{\text{in}}(\varepsilon) = C_0^2(\varepsilon) \frac{A}{V} \quad (3)$$
$$C_0^2(\varepsilon) = |f_0(\varepsilon)|^{-2} = \frac{2\pi/V}{e^{2\pi/V} - 1}$$

is the Gamow factor), which well describes a great amount of experimental data on fusion-reaction cross sections of the type $X(a,b)Y$ in the range 20–100 keV of the colliding energy ε and is used for the extrapolation of $\sigma_0^{\text{in}}(\varepsilon)$ in the limit $\varepsilon \rightarrow 0$

Simple potential model for $X(a,b)Y$

V.S. Melezhik, Nucl.Phys. A550, 223 (1992)

$$\sigma_J^{\text{in}}(\varepsilon) = |f_J(\varepsilon)|^{-2} \frac{A}{V}$$

through the Jost function $f_J(\varepsilon) = |f_J(\varepsilon)| \exp(-i\delta_J)$ of the system $a + X$ determined from the relation

$$k^J |f_J(\varepsilon)|^{-1} R^J \xleftarrow{R \rightarrow 0} \psi_\varepsilon^J(R) \xrightarrow{R \rightarrow \infty} \frac{\sin[kR - \frac{1}{2}J\pi + \delta_J(\varepsilon)]}{kR}$$

and coincident with $|\psi_\varepsilon(R=0)| \exp(-i\delta_J)$ as $J=0$.

L.N. Bogdanova, V.E. Markushin, V.S. Melezhik, L.I. Ponomarev. Sov.J. Nucl. Phys. 34, 662 (1981);
Phys. Lett. B115, 171 (1982)

$$t + d \rightarrow {}^4\text{He} + n \quad J=0$$

two-channel scattering problem

generalized optical potential

$$d + d \rightarrow {}^3\text{He} + n, t + p \quad J=1$$

one-channel problem with nonlocal energy-dependent potential

Simple potential model for X(a,b)Y

V.S. Melezhik, Nucl.Phys. A550, 223 (1992)

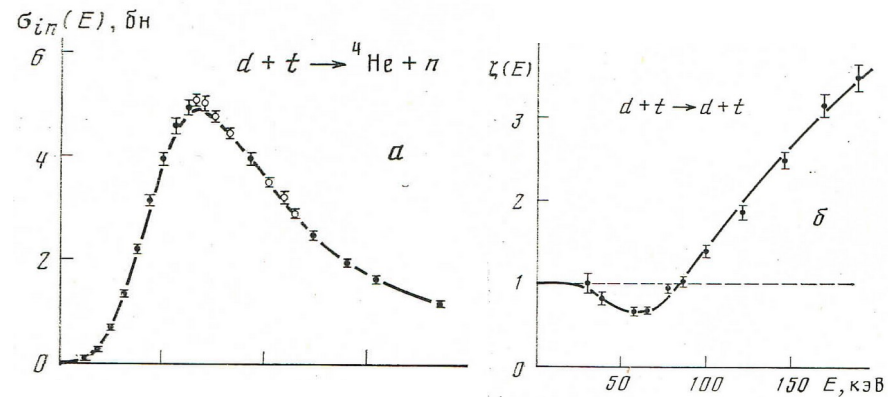
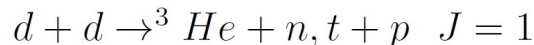
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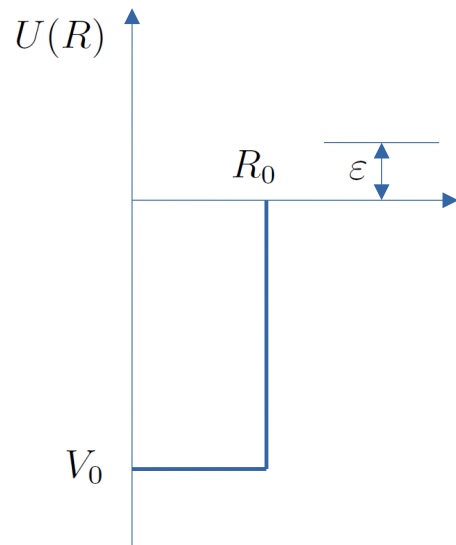
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Simple potential model for X(a,b)Y

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$$U(R) = \begin{cases} -V_0 & R \leq R_0 \\ 0 & R > R_0, \end{cases}$$

$$\sigma_J^{\text{in}}(\varepsilon) = |f_J(\varepsilon)|^{-2} \frac{A}{V}$$

$$\frac{k}{q} |f_0(\varepsilon)|^{-1} \sin qR_0 = \sin(kR_0 + \delta_0), \quad k^2 = 2M\varepsilon$$

$$|f_0(\varepsilon)|^{-1} \cos qR_0 = \cos(kR_0 + \delta_0) \quad q^2 = 2M(\varepsilon + V)$$

$$|f_0(\varepsilon)|^{-2} = \frac{\varepsilon + V_0}{\varepsilon + V_0 \cos^2 qR_0}$$

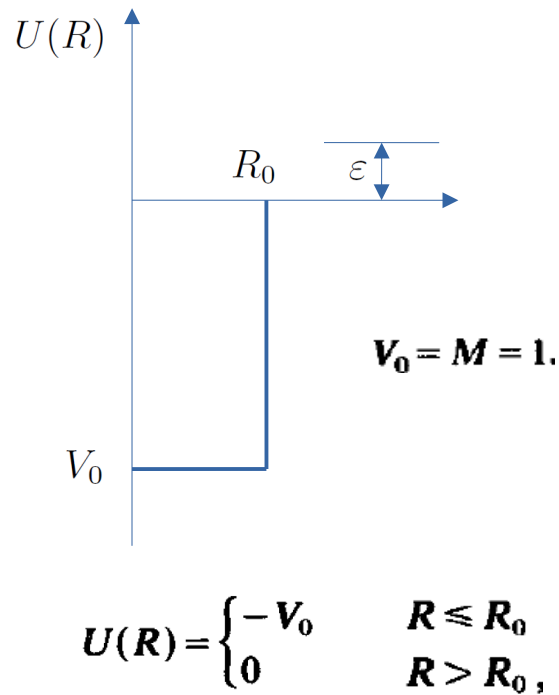
$$|f_0(\varepsilon)|^{-2} = \begin{cases} 1 + \frac{V_0}{\varepsilon} & \text{max} \\ 1 & \text{min} \end{cases}$$

$$n = \nu, \nu + 1, \dots \quad \varepsilon_n = \frac{\pi^2(1 + 2n)^2}{8MR_0^2} - V_0 \quad \varepsilon_{n+1} - \varepsilon_n = \frac{(n + 1)\pi^2}{MR_0^2}$$

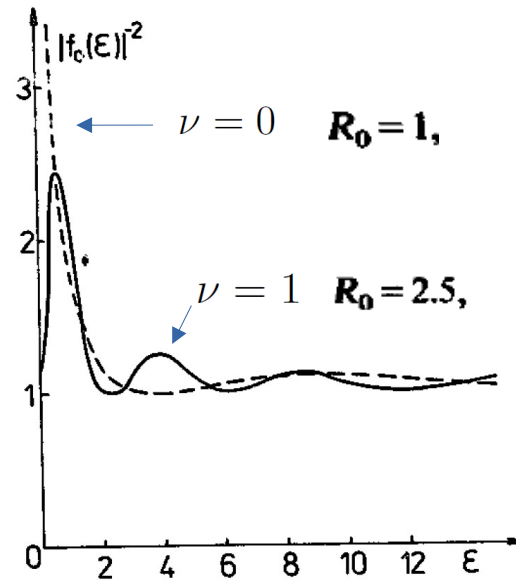
$$\frac{\varepsilon_{n+2} - \varepsilon_{n+1}}{\varepsilon_{n+1} - \varepsilon_n} = \frac{n + 2}{n + 1}$$

Simple potential model for X(a,b)Y

V.S. Melezhik, Nucl.Phys. A550, 223 (1992)



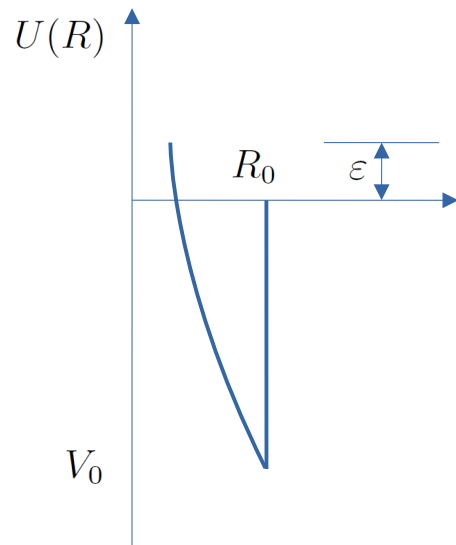
$$\sigma_J^{\text{in}}(\varepsilon) = |f_J(\varepsilon)|^{-2} \frac{A}{V}$$



$$\nu = 1 \quad \frac{\varepsilon_3 - \varepsilon_2}{\varepsilon_2 - \varepsilon_1} = \frac{\nu + 2}{\nu + 1} = \frac{3}{2} \rightarrow \frac{8.66 - 3.93}{3.93 - 0.78} = \frac{4.73}{3.15} = 1.50$$

Simple potential model for X(a,b)Y

V.S. Melezhik, Nucl.Phys. A550, 223 (1992)



$$U(R) = \begin{cases} (Z_1 Z_2 / R) - V_0 & R \leq R_0 \\ 0 & R > R_0 \end{cases}$$

$$\sigma_J^{\text{in}}(\varepsilon) = |f_J(\varepsilon)|^{-2} \frac{A}{V}$$

$$\mathbf{D}(\mathbf{d}, \mathbf{p})\mathbf{T}$$

$$R_0 \sim a_B = 0.525 \times 10^{-8} \text{ cm}$$

$$\mathbf{p}\mu + \mathbf{p} \rightarrow \mathbf{d} + \mathbf{e}^+ + \mu \quad R_0 \sim a_B / 206$$

$$|f_0(\varepsilon)|^{-2} = \frac{C_0^2(q)(\varepsilon + V_0)}{\varepsilon + [V_0 - R_0^{-1} + (\varepsilon + V_0)^{-1} R_0^{-2}] \cos^2 \beta},$$

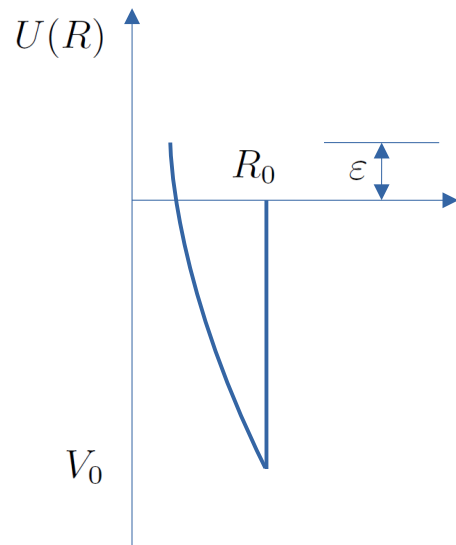
$$\beta = qR_0 - \frac{M}{q} \ln 2qR_0 + \arg \Gamma\left(1 + i\frac{M}{q}\right), \quad q = \sqrt{2M(\varepsilon + V_0)}$$

$$f_{\max}^{-2} = C_0^2(q_n) \left(1 + \frac{V_0}{\varepsilon_n}\right),$$

$$q_n R_0 - \frac{M}{q_n} \ln 2q_n R_0 + \arg \Gamma\left(1 + i\frac{M}{q_n}\right) = \left(\frac{1}{2} + n\right)\pi$$

Simple potential model for X(a,b)Y

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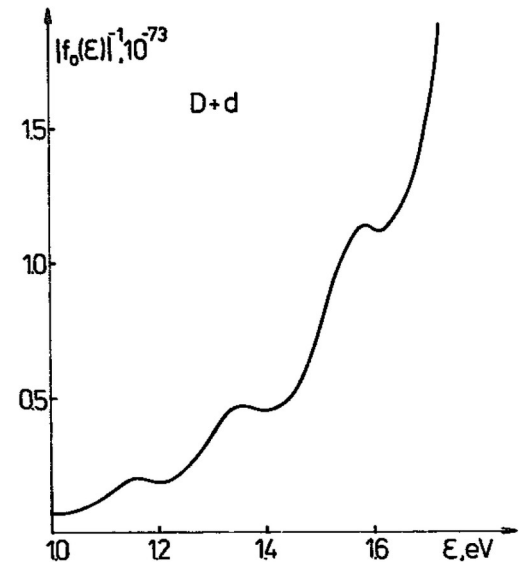
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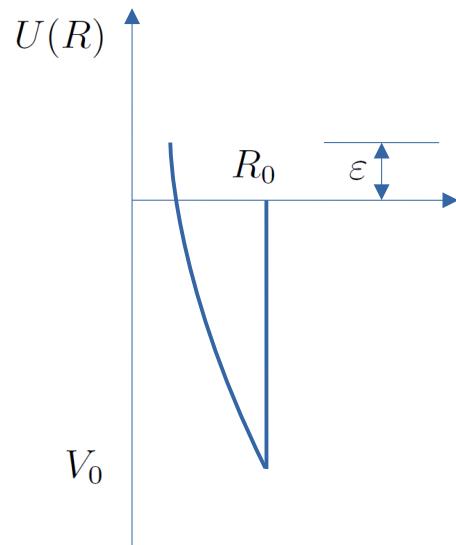
$$f_{\text{max}}^{-2} = C_0^2(q_n) \left(1 + \frac{V_0}{\varepsilon_n}\right)$$



$$\varepsilon_{n+1} - \varepsilon_n \underset{n=\nu}{\cong} \frac{\pi^2(\nu+1)}{MR_0^2} \cong \frac{\pi^2 \times 30}{2 \times 10^3 \times 10} \times 27 \text{ eV} \cong 0.4 \text{ eV}$$

Simple potential model for X(a,b)Y

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$$\sigma_J^{\text{in}}(\epsilon) = |f_J(\epsilon)|^{-2} \frac{A}{V}$$

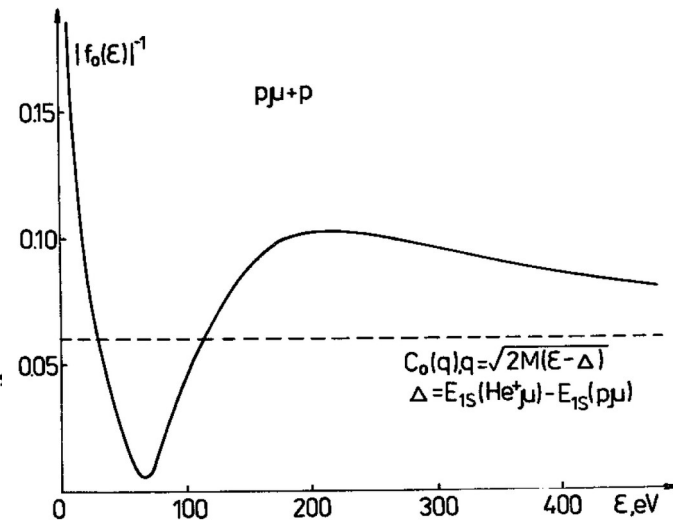
$$p\mu + p \rightarrow d + e^+ + \mu \quad R_0 \sim a_B/206$$

$$f_{\text{max}}^{-2} = C_0^2(q_n) \left(1 + \frac{V_0}{\epsilon_n} \right),$$

$$q_n R_0 - \frac{M}{q_n} \ln 2q_n R_0 + \arg \Gamma \left(1 + i \frac{M}{q_n} \right) = \left(\frac{1}{2} + n \right) \pi$$

$$U(R) = \begin{cases} (Z_1 Z_2 / R) - V_0 & R \leq R_0 \\ 0 & R > R_0 \end{cases}$$

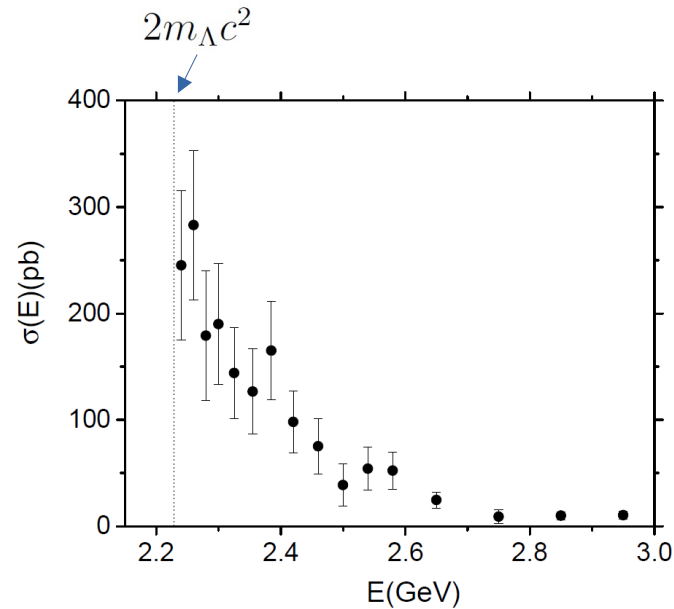
$$\epsilon_{\nu+1} - \epsilon_{\nu} \simeq \frac{\pi^2(\nu+1)}{MR_0^2} \simeq \frac{10 \times 3}{10 \times 10} \times 5 \times 10^3 \simeq 10^2 - 10^3 \text{ eV}$$



Simple potential model for

$$e^+ + e^- \rightarrow \Lambda + \bar{\Lambda}$$

M. Ablikim et. al (BESIII Collaboration) PHYSICAL REVIEW D **107**, 072005 (2023)



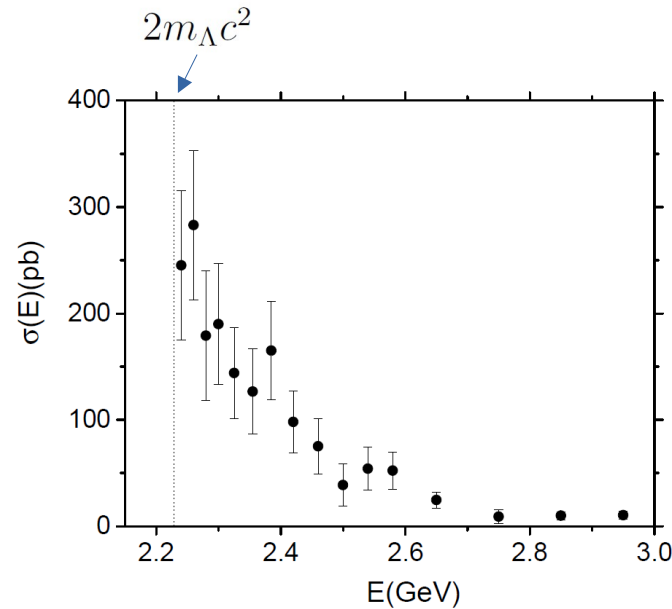
detailed balance principle:

$$\sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) = \frac{k_{\Lambda}^2}{k_e^2} \sigma(\Lambda\bar{\Lambda} \rightarrow e^+e^-) = \frac{k_{\Lambda}^2}{k_e^2} \frac{A}{v_{\Lambda}} |f_J(E)|^{-2} = \frac{k_{\Lambda}}{k_e^2} A |f_J(E)|^{-2}$$

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detailed balance principle:

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$$S^{11}(E) = e^{2i\delta(E)} \frac{1 - 2m_1 k \mathbf{A} |f(E)|^{-2} + iF(E)}{1 + 2m_1 k \mathbf{A} |f(E)|^{-2} + iF(E)}$$

$$\mathbf{A} = \beta |\langle \varphi_E | \xi \rangle|^2, \quad F(E) = \frac{(2m_1)^{1/2}}{\pi} \int_0^{\infty} \frac{\beta |\langle \varphi_E | \xi \rangle|^2 |f(\varepsilon)|^{-2}}{E - \varepsilon} \bar{V}_{\varepsilon} d\varepsilon$$

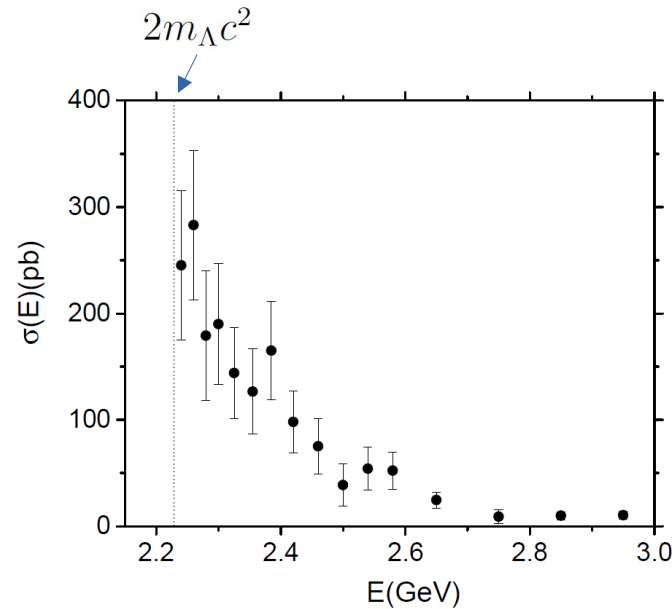
$$V_a = -i\beta |\xi\rangle \langle \xi|$$

$$V_a = -i(2m_2)^{1/2} (E + \Delta)^{1/2} V_{12} |E + \Delta\rangle \langle E + \Delta| V_{21}$$

Simple potential model for

$$e^+ + e^- \rightarrow \Lambda + \bar{\Lambda}$$

M. Ablikim et. al (BESIII Collaboration) PHYSICAL REVIEW D **107**, 072005 (2023)



detailed balance principle:

$$\sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) = \frac{k_{\Lambda}^2}{k_e^2} \sigma(\Lambda\bar{\Lambda} \rightarrow e^+e^-) = \frac{k_{\Lambda}^2}{k_e^2} \frac{A}{v_{\Lambda}} |f_J(E)|^{-2} = \frac{k_{\Lambda}}{k_e^2} A |f_J(E)|^{-2} = \frac{k_{\Lambda}}{k_e^2} A F_D^2(k_e^2) |f_J(E)|^{-2}$$

$$S^{11}(E) = e^{2i\delta(E)} \frac{1 - 2m_1 k \Delta |f(E)|^{-2} + iF(E)}{1 + 2m_1 k \Delta |f(E)|^{-2} + iF(E)}$$

$$A = A' F_D^2(k_e^2) \quad \leftarrow \quad \Delta = \beta |\langle \varphi_E | \xi \rangle|^2, \quad F(E) = \frac{(2m_1)^{1/2}}{\pi} \int_0^{\infty} \frac{\beta |\langle \varphi_E | \xi \rangle|^2 |f(\varepsilon)|^{-2}}{E - \varepsilon} \bar{V}_{\varepsilon} d\varepsilon$$

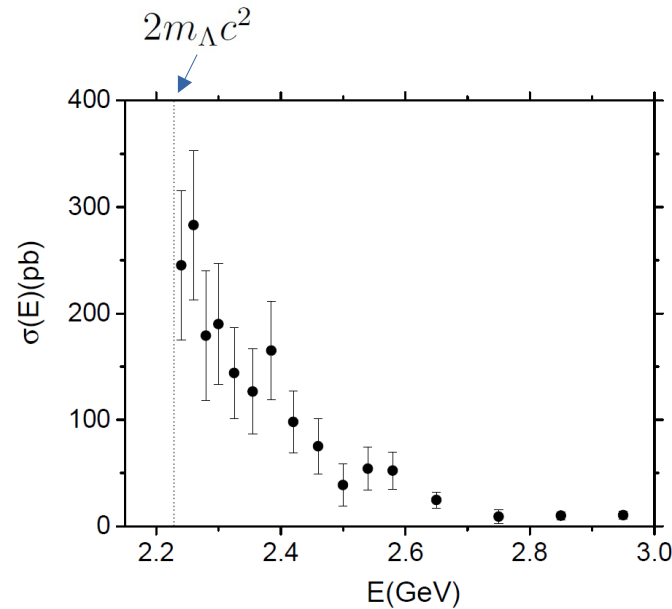
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Simple potential model for

$$e^+ + e^- \rightarrow \Lambda + \bar{\Lambda}$$

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detailed balance principle:

$$\sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) = \frac{k_{\Lambda}^2}{k_e^2} \sigma(\Lambda\bar{\Lambda} \rightarrow e^+e^-) = \frac{k_{\Lambda}^2}{k_e^2} \frac{A}{v_{\Lambda}} |f_J(E)|^{-2} = \frac{k_{\Lambda}}{k_e^2} A |f_J(E)|^{-2} = \frac{k_{\Lambda}}{k_e^2} A F_D^2(k_e^2) |f_J(E)|^{-2}$$

$\Lambda\bar{\Lambda}$ interaction in final state (S-wave):

$$\sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) = \frac{k_{\Lambda}}{k_e^2} F_D^2(k_e^2) g |\psi(R=0)|^2$$

$$\text{dipole formfactor } F_D = (1 - \frac{k_e^4}{\Delta^2})^{-2} \quad \Delta \sim 1\text{GeV}$$

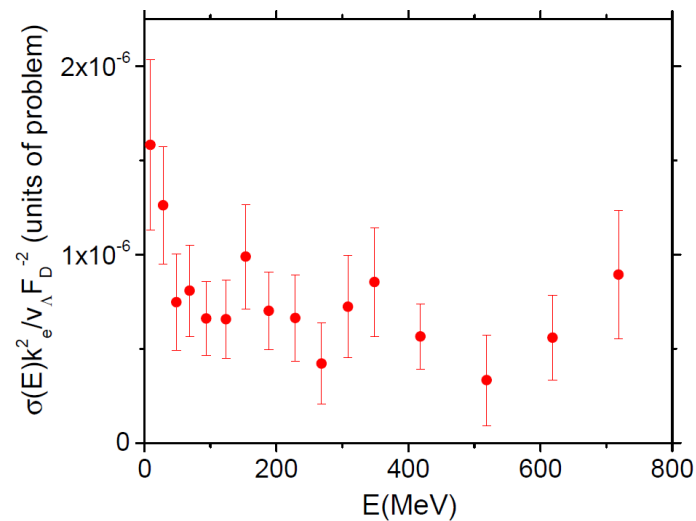
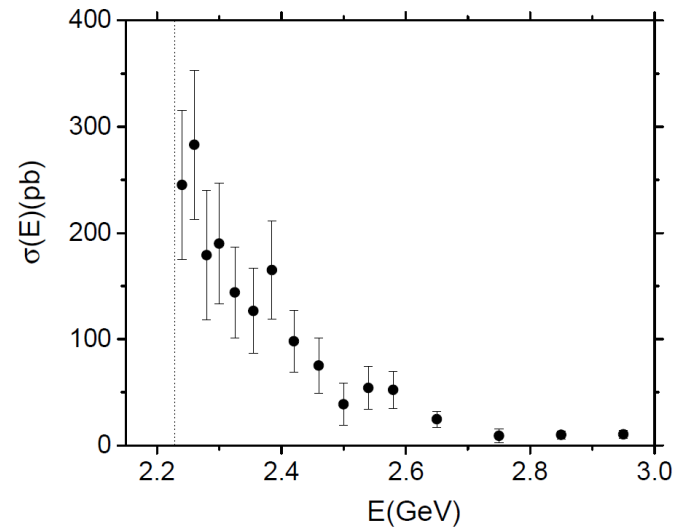
A.I.Milstein, S.G. Salnikov, JETP Letters 117 (2023); Phys Rev D106 (2022)

J. Haidenbauer, X-G. Kang, U-G. Meissner, Nucl Phys A929 (2016)

Simple potential model for

$$e^+ + e^- \rightarrow \Lambda + \bar{\Lambda}$$

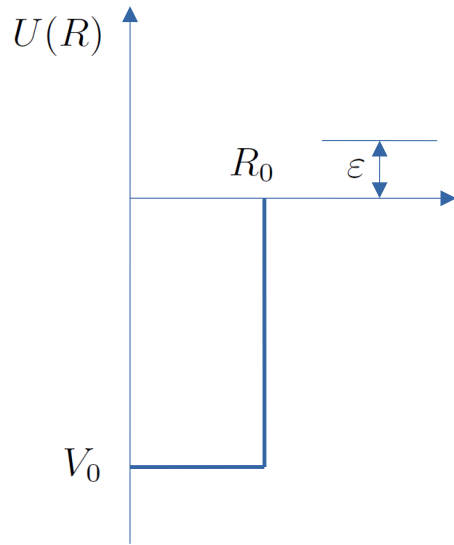
$$A|f_J(E)|^{-2} = \sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) \frac{k_e^2}{k_\Lambda F_D^2}$$



$$\Delta = 1.8\text{GeV}$$

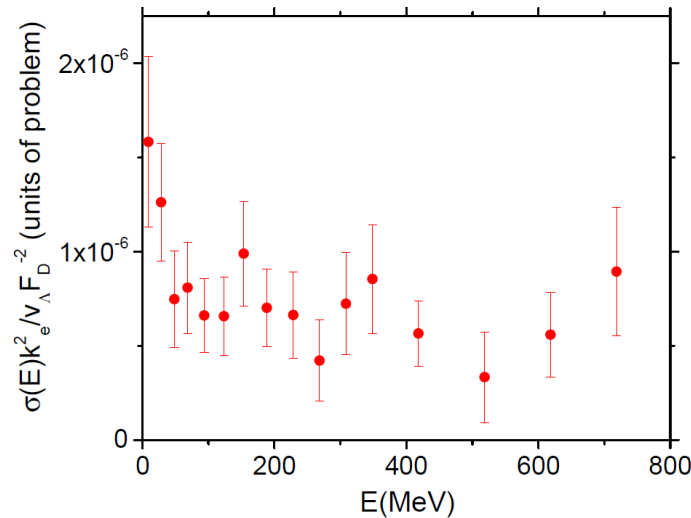
Simple potential model for

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$$U(R) = \begin{cases} -V_0 & R \leq R_0 \\ 0 & R > R_0, \end{cases}$$

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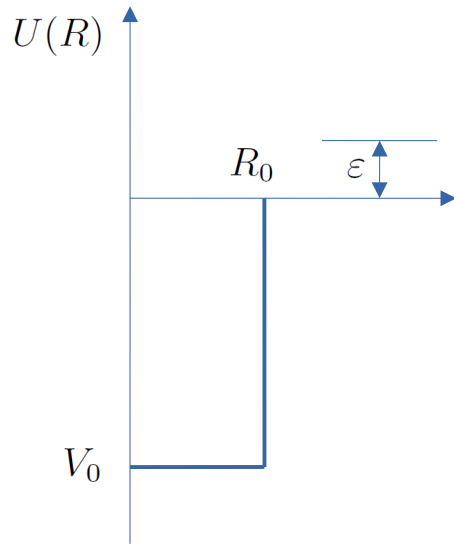


$$\epsilon_n = \frac{\pi^2(1+2n)^2}{8MR_0^2} - V_0 \quad n = \nu, \nu+1, \dots$$

$$\epsilon_{n+1} - \epsilon_n = \frac{\pi^2(n+1)}{MR_0^2}$$

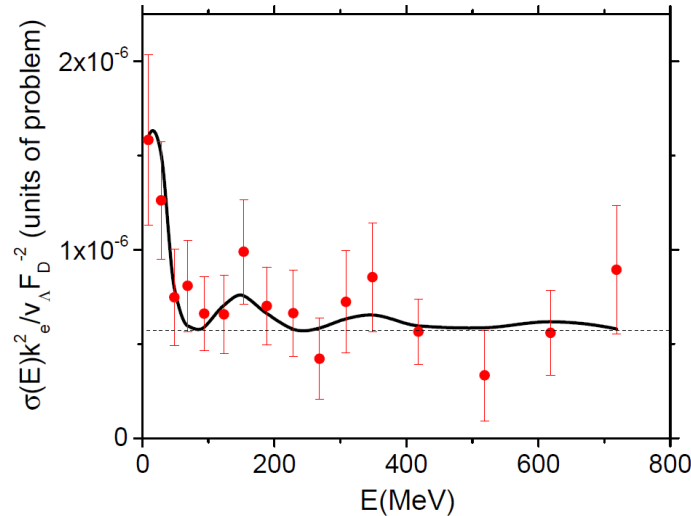
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$$A|f_J(E)|^{-2} = \sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) \frac{k_e^2}{k_\Lambda F_D^2}$$



$$V_0 = 72 \text{ MeV} \quad R_0 = 3.26 \text{ fm}$$

$$A = 0.57 \times 10^{-6}$$

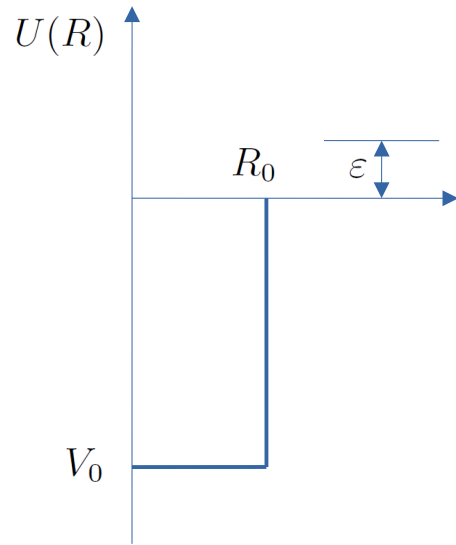
$$\chi^2 = 12/16$$

$$\epsilon_n = \frac{\pi^2(1+2n)^2}{8MR_0^2} - V_0 \quad n = \nu, \nu+1, \dots$$

$$\epsilon_{n+1} - \epsilon_n = \frac{\pi^2(n+1)}{MR_0^2}$$

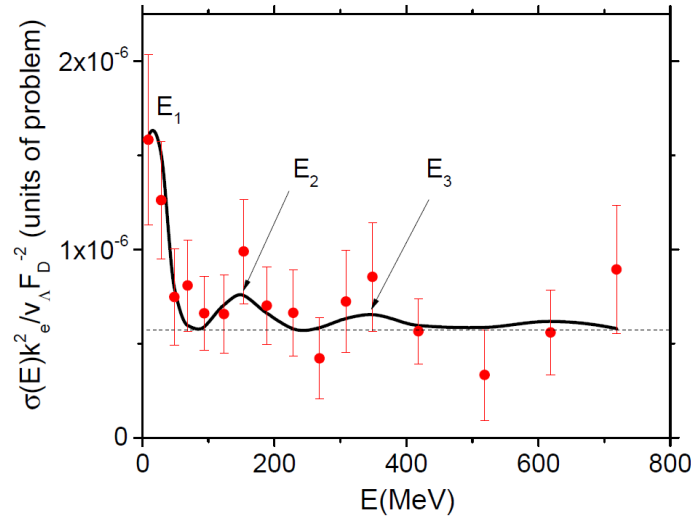
Simple potential model for

$$e^+ + e^- \rightarrow \Lambda + \bar{\Lambda}$$



$$U(R) = \begin{cases} -V_0 & R \leq R_0 \\ 0 & R > R_0 \end{cases}$$

$$A|f_J(E)|^{-2} = \sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) \frac{k_e^2}{k_\Lambda F_D^2}$$



$$V_0 = 72 \text{ MeV} \quad R_0 = 3.26 \text{ fm}$$

$$A = 0.57 \times 10^{-6}$$

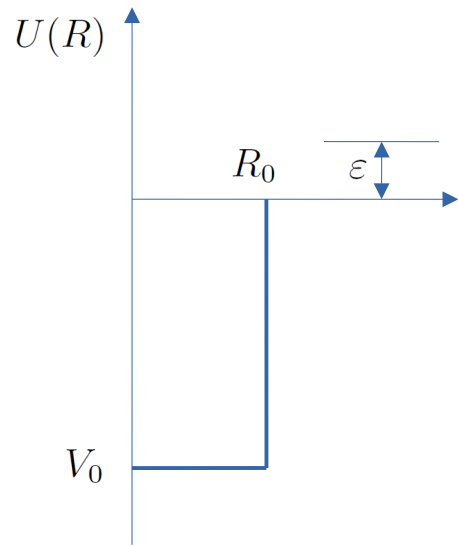
$$\chi^2 = 7/16$$

$$\nu = 1 \quad \frac{\nu+2}{\nu+1} = 1.5$$

$$\frac{E_3 - E_2}{E_2 - E_1} = 1.4$$

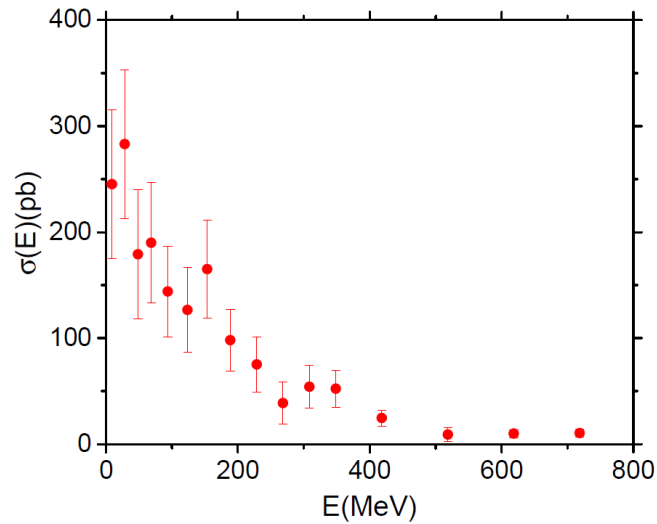
Simple potential model for

$$e^+ + e^- \rightarrow \Lambda + \bar{\Lambda}$$



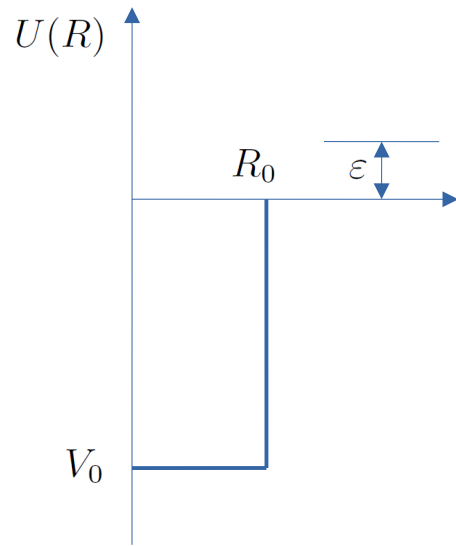
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$$\sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) = \frac{k_\Lambda}{k_e^2} F_D^2 A |f_0(E)|^{-2}$$



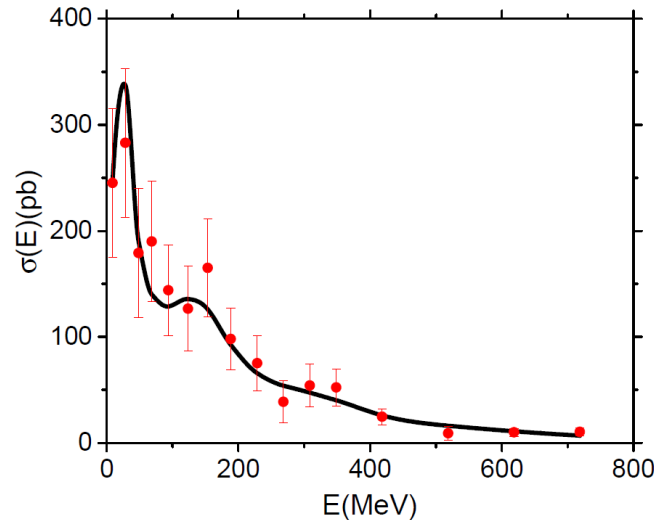
Simple potential model for

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$$\Delta = 1.8\text{GeV} \quad A = 0.57 \times 10^{-6} \quad V_0 = 72\text{MeV} \quad R_0 = 3.26\text{fm}$$

$$\chi^2 = 7/16$$

Conclusion & Outlook

- simple potential model $\sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) = \frac{k_\Lambda}{k_e^2} F_D^2 A |f_0(E)|^{-2}$

fitting parameters: $A \quad V_0 \quad R_0$

Conclusion & Outlook

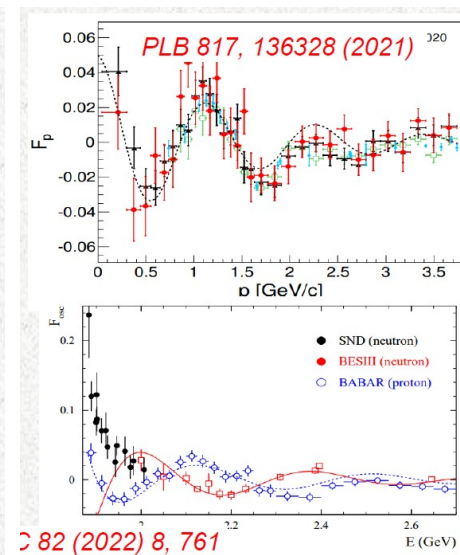
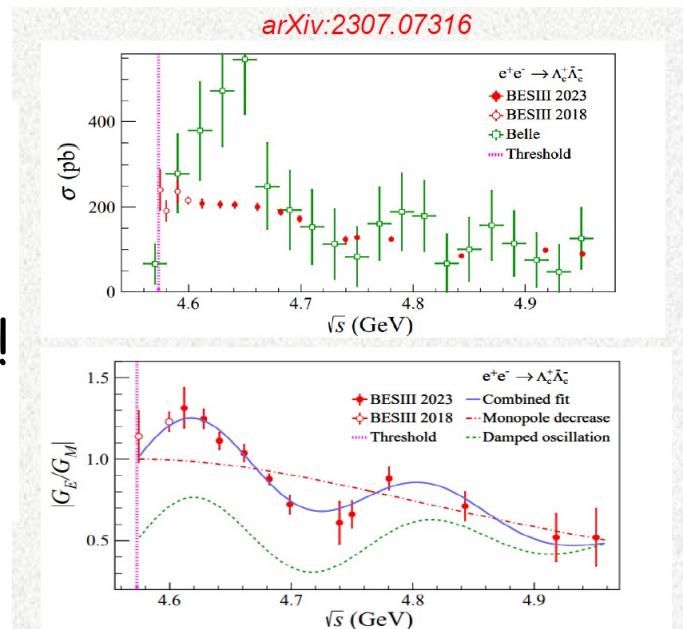
- simple potential model

$$\sigma(e^+e^- \rightarrow \Lambda\bar{\Lambda}) = \frac{k_\Lambda}{k_e^2} F_D^2 A |f_0(E)|^{-2}$$

fitting parameters: A V_0 R_0

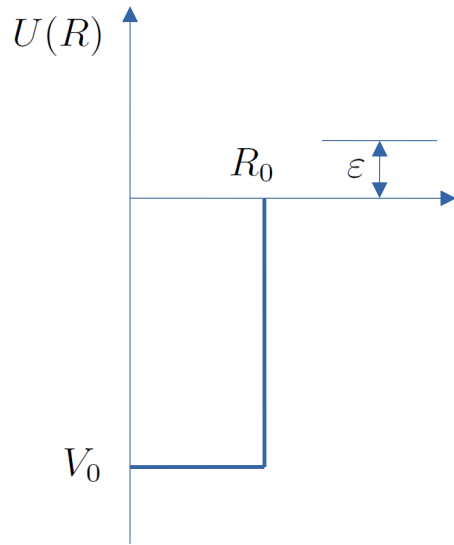
- $e^+e^- \rightarrow \Lambda_c\bar{\Lambda}_c$
formfactors ?!

S, D-waves
Coulomb



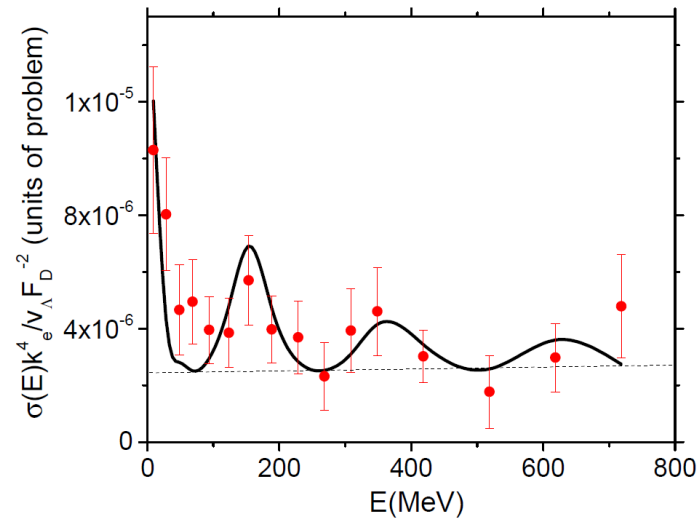
Simple potential model for

$$e^+ + e^- \rightarrow \Lambda + \bar{\Lambda}$$



$$U(R) = \begin{cases} -V_0 & R \leq R_0 \\ 0 & R > R_0 \end{cases}$$

$$\sigma_J^{\text{in}}(\epsilon) = |f_J(\epsilon)|^{-2} \frac{A}{V}$$



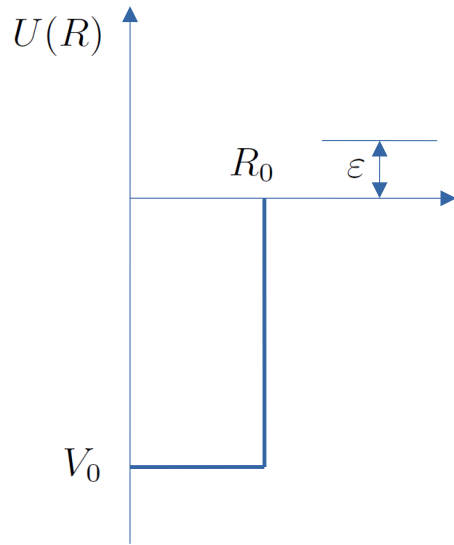
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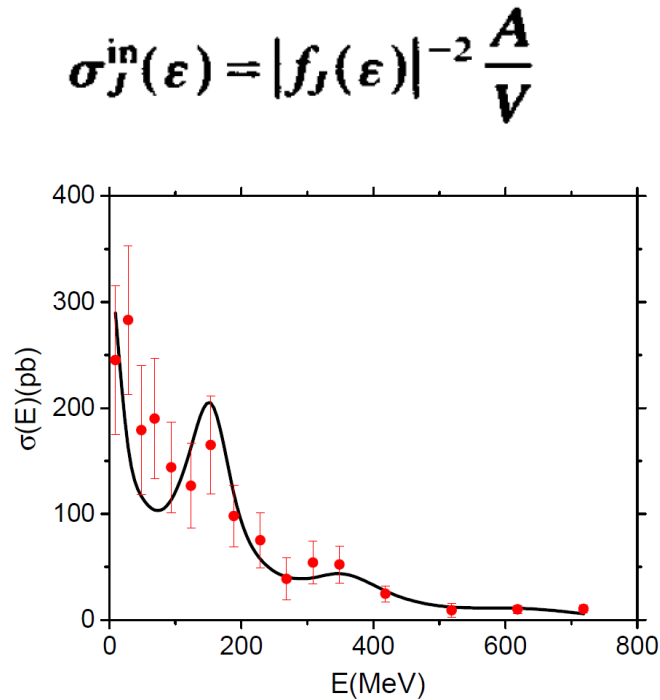
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Simple potential model for

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